FFTs in Graphics and Vision

Spherical Harmonics
Outline

Math Stuff

Review

Finding the Spherical Harmonics
Homogenous Polynomials

A homogenous polynomial of degree $d$ in $n$ variables can be expressed as:

$$ p_d(x_1, \cdots, x_n) = \sum_{j_1+\cdots+j_n=d} a_{j_1\cdots j_n} \cdot x_1^{j_1} \cdots x_n^{j_n} $$

$$ = \sum_{j_1=0}^{d} x_1^{j_1} \left( \sum_{j_2+\cdots+j_n=d-j_1} a_{j_1\cdots j_n} \cdot x_2^{j_2} \cdots x_n^{j_n} \right) $$
Homogenous Polynomials

\[ p_d(x_1, \ldots, x_n) = \sum_{j_1=0}^{d} x_1^{j_1} \left( \sum_{j_2+\ldots+j_n=d-j_1} a_{j_1\ldots j_n} \cdot x_2^{j_2} \cdots x_n^{j_n} \right) \]

If we fix the value of the first coefficient at \( x_1 = \zeta \), we get a new polynomial in \( n - 1 \) variables:

\[ q_d(x_2, \ldots, x_n) = p_d(\zeta, x_2, \ldots, x_n) \]

This gives:

\[ q_d(x_2, \ldots, x_n) = \sum_{j_1=0}^{d} \zeta^{j_1} \left( \sum_{j_2+\ldots+j_n=d-j_1} a_{j_1\ldots j_n} \cdot x_2^{j_2} \cdots x_n^{j_n} \right) \]
Homogenous Polynomials

\[ q_d(x_2, \cdots, x_n) = \sum_{j_1=0}^{d} \zeta^{j_1} \left( \sum_{j_2+\cdots+j_n = d-j_1} a_{j_1 \cdots j_n} \cdot x_2^{j_2} \cdots x_n^{j_n} \right) \]

The new polynomial, obtained by fixing the value of the first variable, is a polynomial of degree at most \( d \) in \( n - 1 \) variables.
Review

So far, we have considered the representation of the 2D group of rotations, acting on the space of (complex-valued) functions on the unit circle:

\[(\rho_R f)(\rho) = f(R^{-1}\rho)\]
Since the group of 2D rotations is commutative, Schur’s lemma tells us that the space of functions can be expressed as the sum of one-dimensional irreducible representations:

\[
V = \bigoplus V^k \quad \text{w/} \quad \dim(V^k) = 1 \\
\rho_R(V^k) = V^k \quad \forall R \in SO(2)
\]
Review

In the 2D case, we know that the $V^k$ are spanned by the complex exponentials of frequency $k$:

$$V^k = \{a \cdot e^{i k \theta} \mid a \in \mathbb{C}\}$$

$\Rightarrow$ Computing the Fourier transform of a function:

$$f(\theta) = \sum_{k=-\infty}^{\infty} \hat{f}_k \cdot e^{i k \theta}$$

enables fast smoothing/correlation by allowing us to only consider inner-products within $V^k$. 
Functions on the Sphere

What happens when we consider the space of functions on the unit sphere?

Since the group of 3D rotations doesn’t commute, the irreducible representations:

\[ V = \bigoplus V^k \]

do not have to be one-dimensional.
Functions on the Sphere

What happens when we consider the space of functions on the unit sphere?

However, we would still like to compute the irreducible representations. And in particular…
Goal

Let $F$ be the space of (complex-value) functions on the unit sphere and let $\rho$ be the representation of the group of 3D rotations, acting on the space of functions by rotation:

$$(\rho_R f)(p) = f(R^{-1}p)$$
Goal

Let $F$ be the space of (complex-value) functions on the unit sphere and let $\rho$ be the representation of the group of 3D rotations, acting on the space of functions by rotation:

$$(\rho_R f)(p) = f(R^{-1}p)$$

What are the irreducible representations?
What We Know

We know that the irreducible representations are related to the sub-spaces of homogenous polynomials of fixed degree:

- Given a homogenous polynomial of degree $l$, a rotation of the polynomial will still be a homogenous polynomial of degree $l$:
  \[
  \rho_R \left( HP^l(x, y, z) \right) = HP^l(x, y, z)
  \]

- If we “throw out” the homogenous polynomials whose restriction to the unit sphere can be expressed as the restriction of a homogenous polynomial of smaller degree, we get a $(2l + 1)$-dimensional representation.
Letting $l$ index the degree of the homogenous polynomial, we can decompose the space of spherical functions as:

$$V = \bigoplus_{l=0}^{\infty} V^l \text{ w/ } \dim(F^l) = 2l + 1$$

$$\rho_R(V^l) = V^l \quad \forall R \in SO(3)$$

with $f \in V^l$ expressable as the restriction of a homogenous polynomial in three variables, of degree $l$ (and orthogonal to all homogenous polynomials of degree $l - 2$).
What We Want to Know

What is a basis for each $V^l$?

$$V^l = \text{Span}\{f^0_l, \ldots, f^{2l}_l\}$$
Parameterizing Spherical Functions

A point on the unit sphere can be parameterized by its angles of longitude and latitude:

\[ \Phi(\theta, \phi) = (\cos \theta \cdot \sin \phi, \cos \phi, \sin \theta \cdot \sin \phi) \]

with \( \theta \in [0, 2\pi) \) and \( \phi \in [0, \pi] \).
Parameterizing Spherical Functions

\[ \Phi(\theta, \phi) = (\cos \theta \cdot \sin \phi, \cos \phi, \sin \theta \cdot \sin \phi) \]

Fixing \( \theta \), we get great semi-circles (meridians) through the North and South poles.
Parameterizing Spherical Functions

\[ \Phi(\theta, \phi) = (\cos \theta \cdot \sin \phi, \cos \phi, \sin \theta \cdot \sin \phi) \]

Fixing \( \phi \), we get parallels about the \( y \)-axis.
Parameterizing Spherical Functions

\[ \Phi(\theta, \phi) = (\cos \theta \cdot \sin \phi, \cos \phi, \sin \theta \cdot \sin \phi) \]

A spherical function can be represented by its values at every spherical angle.
Rotations About the $\gamma$-Axis

Instead of considering the action of the entire group of rotations on the space of spherical functions, we can consider the subset of rotations that rotate about the $\gamma$-axis:

$$(\rho_{R\gamma}(\alpha)f)(\theta, \phi) = f(\theta - \alpha, \phi)$$
Rotations About the $y$-Axis

Instead of considering the entire group of rotations, we can consider the subset of rotations that rotate about the $y$-axis.

- This set of rotations is a group:
  - The product of two rotations about the $y$-axis is still a rotation about the $y$-axis.
  - Rotation by $-\alpha$ degrees about the $y$-axis is the inverse of the rotation by $\alpha$ degrees.
Rotations About the $y$-Axis

Instead of considering the entire group of rotations, we can consider the subset of rotations that rotate about the $y$-axis.

- This set of rotations is a group.
- This sub-group is commutative.
Rotations About the $y$-Axis

⇒ The sub-spaces $V^l$ are representations for the commutative sub-group of rotations about the $y$-axis.

⇒ Each $V^l$ can be expressed as the sum of one-dimensional representations that are fixed by rotations about the $y$-axis.
Rotations About the $y$-Axis

Thus, for each $l$, there must exist a basis of orthogonal functions $\{f_l^0(\theta, \phi), \ldots, f_l^{2l}(\theta, \phi)\}$ such that a rotation by $\alpha$ degrees about the $y$-axis becomes multiplication by a complex number:

$$\rho_{R_y(\alpha)}(f_l^k) = \chi_l^k(\alpha) \cdot f_l^k$$
Rotations About the $y$-Axis

$$\rho_{R_y}(\alpha)\left(f_i^k\right) = \chi_i^k(\alpha) \cdot f_i^k$$

Since the representation is unitary, we know that for any angle of rotation $\alpha$, we must have:

$$\|\chi_i^k(\alpha)\| = 1$$

Since we know that representations preserve the group structure, and since rotating by $\alpha$ degrees and then by $\beta$ degrees is equivalent to rotating by $(\alpha + \beta)$ degrees, we have:

$$\chi_i^k(\alpha + \beta) = \chi_i^k(\alpha) \cdot \chi_i^k(\beta)$$
Rotations About the $y$-Axis

$$\rho_{R_y}(\alpha)(f_l^k) = \chi^k_l(\alpha) \cdot f_l^k$$

The only functions satisfying these properties are:

$$\chi^k_l(\alpha) = e^{-im_{kl}\alpha}$$

Moreover, since we know that rotations by $\alpha = 2\pi$ degrees about the $y$-axis do not change a function, the powers $m_{kl}$ must be integers.
Rotations About the $y$-Axis

$$\rho_{R_y(\alpha)}(f_l^k) = \chi_l^k(\alpha) \cdot f_l^k$$

Consider the function:

$$\tilde{f}_l^k(\theta, \phi) = \frac{f_l^k(\theta, \phi)}{e^{im_{kl}\theta}}$$
Rotations About the $y$-Axis

$$\tilde{f}_l^k (\theta, \phi) = \frac{f_l^k(\theta, \phi)}{e^{im_{kl} \theta}}$$

When we rotate by $\alpha$ degrees about the $y$-axis:

$$\left( \rho_{R_y(\alpha)} \tilde{f}_l^k \right)(\theta, \phi) = \tilde{f}_l^k(\theta - \alpha, \phi)$$

$$= \frac{f_l^k(\theta - \alpha, \phi)}{e^{im_{kl}(\theta-\alpha)}}$$

$$= e^{-im_{kl} \alpha} \cdot \frac{f_l^k(\theta, \phi)}{e^{im_{kl} \alpha} \cdot e^{im_{kl} \theta}}$$

$$= \tilde{f}_l^k(\theta, \phi)$$
Rotations About the $y$-Axis

$$\tilde{f}_l^k(\theta, \phi) = \frac{f_l^k(\theta, \phi)}{e^{im_{kl}\theta}}$$

When we rotate by $\alpha$ degrees about the $y$-axis:

$$\left( \rho_{R_y(\alpha)} \tilde{f}_l^k \right) = \tilde{f}_l^k$$

Since these functions are unchanged by rotations about the $y$-axis, they are only functions of $\phi$:

$$\tilde{f}_l^k(\theta, \phi) = g_l^k(\phi)$$

So we have:

$$f_l^k(\theta, \phi) = e^{im_{kl}\theta} \cdot g_l^k(\phi)$$
Factoring the Spherical Harmonics

\[ f^k_l(\theta, \phi) = e^{im_k\theta} \cdot g^k_l(\phi) \]

What can we say about the integers \( m_{kl} \)?

The \((x, y, z)\) coordinates of a point on the sphere are defined by:

\[ \Phi(\theta, \phi) = (\cos \theta \cdot \sin \phi, \cos \phi, \sin \theta \cdot \sin \phi) \]

\[ \downarrow \]

\[ \cos \theta = \frac{x}{\sin \phi} \quad \cos \phi = y \quad \sin \theta = \frac{z}{\sin \phi} \]
Factoring the Spherical Harmonics

\[ f^k_l(\theta, \phi) = e^{im_{kl}\theta} \cdot g^k_l(\phi) \]

What can we say about the integers \( m_{kl} \)?

Fixing the angle of latitude, \( \phi = \phi_0 \), gives:

\[
\begin{align*}
 f^k_l(\theta, \phi_0) &= (\cos \theta + i \sin \theta)^{m_{kl}} \cdot g^k_l(\phi_0) \\
 &= \left( \frac{x}{\sin \phi_0} + i \frac{z}{\sin \phi_0} \right)^{m_{kl}} \cdot g^k_l(\phi_0) \\
 &= (x + iz)^{m_{kl}} \cdot \frac{g^k_l(\phi_0)}{\sin^{m_{kl}} \phi_0}
\end{align*}
\]
Factoring the Spherical Harmonics

\[ f^k_l(\theta, \phi) = e^{im_{kl}\theta} \cdot g^k_l(\phi) \]

What can we say about the integers \( m_{kl} \)?

Fixing the angle of latitude, \( \phi = \phi_0 \), gives:

\[ f^k_l(\theta, \phi_0) = (x + iz)^{m_{kl}} \cdot \frac{g^k_l(\phi_0)}{\sin^{m_{kl}}\phi_0} \]

But \( f^k_l(\theta, \phi) \) is the restriction of a homogenous polynomial of degree \( l \) to the unit sphere.

So fixing \( y = \cos \phi_0 \), we get a polynomial of degree at most \( l \):

\[ -l \leq m_{kl} \leq l \]
The Spherical Harmonics

Summarizing:

We know that the space of spherical functions $V$ can be expressed as the sum of sub-representations:

$$V = \bigoplus_{l=0}^{\infty} V^l \quad \text{w/} \quad \dim(V^l) = 2l + 1$$

$$\rho_R(V^l) = V^l \quad \forall R \in SO(3)$$

where the functions in $V^l$ are obtained by considering the restrictions of homogenous polynomials of degree $l$ to the unit sphere.
The Spherical Harmonics

Summarizing:

Each $V_l$ is spanned by an orthogonal basis
$\{f_l^0(\theta, \phi), \ldots, f_l^{2l}(\theta, \phi)\}$ where the $k$-th basis function can be expressed as:

$$f_l^k(\theta, \phi) = e^{im_{kl}\theta} \cdot g_l^k(\phi)$$

where $m_{kl}$ is an integer in the range $[-l, l]$. 
The Spherical Harmonics

Summarizing:

It turns out that for every value of $-l \leq m \leq l$ there is exactly one basis function:

$$Y_l^m(\theta, \phi) = e^{im\theta} \cdot g_l^m(\phi)$$

These are the spherical harmonics of degree $l$. 
Aside

To evaluate the spherical harmonics, we need to know what the functions $g^m_l(\phi)$ are.

These are defined by setting:

$$g^m_l(\phi) = P_l^{|m|}(\cos \phi)$$

where the $P_l^m$ are the associated Legendre polynomials, defined by:

$$P_l^m(x) = \frac{(-1)^m}{2^l l!} (1 - x^2)^{m/2} \frac{d^{l+m}}{dx^{l+m}} (m^2 - 1)^l.$$
The Spherical Harmonics

Examples \((l = 0)\):

\[ Y_0^0(\theta, \phi) = \frac{1}{\sqrt{4\pi}} \]
The Spherical Harmonics

Examples \((l = 1)\):

\[
Y_{1}^{-1}(\theta, \phi) = \sqrt{\frac{3}{8\pi}} \sin \phi \cdot e^{-i\theta}
\]

\[
Y_{1}^{0}(\theta, \phi) = \sqrt{\frac{3}{4\pi}} \cos \phi
\]

\[
Y_{1}^{1}(\theta, \phi) = \sqrt{\frac{3}{8\pi}} \sin \phi \cdot e^{i\theta}
\]
The Spherical Harmonics

Examples \((l = 2)\):

\[
\begin{align*}
Y_{2}^{-2}(\theta, \phi) &= \sqrt{\frac{15}{32\pi}} \sin^2 \phi \cdot e^{-i2\theta} \\
Y_{2}^{-1}(\theta, \phi) &= \sqrt{\frac{15}{8\pi}} \sin \phi \cdot \cos \phi \cdot e^{-i\theta} \\
Y_{2}^{0}(\theta, \phi) &= \sqrt{\frac{5}{16\pi}} (3 \cos^2 \phi - 1) \\
Y_{2}^{1}(\theta, \phi) &= \sqrt{\frac{15}{8\pi}} \sin \phi \cdot \cos \phi \cdot e^{i\theta} \\
Y_{2}^{2}(\theta, \phi) &= \sqrt{\frac{15}{32\pi}} \sin^2 \phi \cdot e^{i2\theta}
\end{align*}
\]
The Spherical Harmonics

\[
\begin{align*}
\cos \theta &= \frac{x}{\sin \phi} & \cos \phi &= y & \sin \theta &= \frac{z}{\sin \phi} \\
Y_0^0(\theta, \phi) &= \frac{1}{\sqrt{4\pi}}
\end{align*}
\]
The Spherical Harmonics

\[
\cos \theta = \frac{x}{\sin \phi} \quad \cos \phi = y \quad \sin \theta = \frac{z}{\sin \phi}
\]

\[
Y_{1}^{-1}(\theta, \phi) = \sqrt{\frac{3}{8\pi}} \sin \phi \cdot e^{-i\theta}
\]

\[
= \sqrt{\frac{3}{8\pi}} \sin \phi \cdot (\cos \theta - i \sin \theta)
\]

\[
= \sqrt{\frac{3}{8\pi}} \sin \phi \cdot \left( \frac{x}{\sin \phi} - i \frac{z}{\sin \phi} \right)
\]

\[
= \sqrt{\frac{3}{8\pi}} (x - iz)
\]
The Spherical Harmonics

\[
\begin{align*}
\cos \theta &= \frac{x}{\sin \phi} \quad \cos \phi = y \quad \sin \theta = \frac{z}{\sin \phi} \\
Y_1^0(\theta, \phi) &= \sqrt{\frac{3}{4\pi}} \cos \phi \\
&= \sqrt{\frac{3}{4\pi}} y
\end{align*}
\]
The Spherical Harmonics

\[
\begin{align*}
\cos \theta &= \frac{x}{\sin \phi} \quad \cos \phi &= y \quad \sin \theta &= \frac{z}{\sin \phi} \\

\mathbf{Y}_1^1(\theta, \phi) &= \sqrt{\frac{3}{8\pi}} \sin \phi \cdot e^{i\theta} \\
&= \sqrt{\frac{3}{8\pi}} \sin \phi \cdot (\cos \theta + i \sin \theta) \\
&= \sqrt{\frac{3}{8\pi}} \sin \phi \cdot \left( \frac{x}{\sin \phi} + i \frac{z}{\sin \phi} \right) \\
&= \sqrt{\frac{3}{8\pi}} (x + iz)
\end{align*}
\]
The Spherical Harmonics

\[
\cos \theta = \frac{x}{\sin \phi} \quad \cos \phi = y \quad \sin \theta = \frac{z}{\sin \phi}
\]

\[
Y_{2}^{-2}(\theta, \phi) = \sqrt{\frac{15}{32\pi}} \sin^2 \phi \cdot e^{-i2\theta}
\]

\[
= \sqrt{\frac{15}{32\pi}} \sin^2 \phi \cdot (\cos \theta - i \sin \theta)^2
\]

\[
= \sqrt{\frac{15}{32\pi}} \sin^2 \phi \cdot \left( \frac{x}{\sin \phi} - i \frac{z}{\sin \phi} \right)^2
\]

\[
= \sqrt{\frac{15}{32\pi}} (x - iz)^2
\]
The Spherical Harmonics

\[
\begin{align*}
\cos \theta &= \frac{x}{\sin \phi} & \cos \phi &= y & \sin \theta &= \frac{z}{\sin \phi} \\
\end{align*}
\]

\[
Y_{2}^{-1}(\theta, \phi) = \sqrt{\frac{15}{8\pi}} \sin \phi \cdot \cos \phi \cdot e^{-i\theta} \\
= \sqrt{\frac{15}{8\pi}} y \cdot \sin \phi \cdot (\cos \theta - i \sin \theta) \\
= \sqrt{\frac{15}{8\pi}} y \cdot \sin \phi \cdot \left( \frac{x}{\sin \phi} - i \frac{z}{\sin \phi} \right) \\
= \sqrt{\frac{15}{8\pi}} y \cdot (x - iz)
\]
The Spherical Harmonics

\[
\begin{align*}
\cos \theta &= \frac{x}{\sin \phi} & \cos \phi &= y & \sin \theta &= \frac{z}{\sin \phi} \\
Y_2^0(\theta, \phi) &= \sqrt{\frac{5}{16\pi}} (3 \cos^2 \phi - 1) \\
&= \sqrt{\frac{5}{16\pi}} (3y^2 - 1) \\
&= \sqrt{\frac{5}{16\pi}} (2y^2 - x^2 - z^2)
\end{align*}
\]
The Spherical Harmonics

\[
\begin{align*}
\cos \theta &= \frac{x}{\sin \phi} & \cos \phi &= y & \sin \theta &= \frac{z}{\sin \phi} \\
\end{align*}
\]

\[
Y^1_2(\theta, \phi) = \sqrt{\frac{15}{8\pi}} \sin \phi \cdot \cos \phi \cdot e^{i\theta} \\
= \sqrt{\frac{15}{8\pi}} y \cdot \sin \phi \cdot (\cos \theta + i \sin \theta) \\
= \sqrt{\frac{15}{8\pi}} y \cdot \sin \phi \cdot \left( \frac{x}{\sin \phi} + i \frac{z}{\sin \phi} \right) \\
= \sqrt{\frac{15}{8\pi}} y \cdot (x + iz)
\]
The Spherical Harmonics

\[ \cos \theta = \frac{x}{\sin \phi} \quad \cos \phi = y \quad \sin \theta = \frac{z}{\sin \phi} \]

\[ Y_2^2(\theta, \phi) = \sqrt{\frac{15}{32\pi}} \sin^2 \phi \cdot e^{i2\theta} \]

\[ = \sqrt{\frac{15}{32\pi}} \sin^2 \phi \cdot (\cos \theta + i \sin \theta)^2 \]

\[ = \sqrt{\frac{15}{32\pi}} \sin^2 \phi \cdot \left( \frac{x}{\sin \phi} + i \frac{z}{\sin \phi} \right)^2 \]

\[ = \sqrt{\frac{15}{32\pi}} (x + iz)^2 \]
Implications

For any spherical function $f(\theta, \phi)$, we can express $f$ as the sum of functions in $V^l$:

$$f(\theta, \phi) = \sum_{l=0}^{\infty} f_l(\theta, \phi)$$

Each $f_l$ can be expressed as the sum of harmonics:

$$f_l(\theta, \phi) = \sum_{k=-l}^{l} \hat{f}_{lk} \cdot Y^k_l(\theta, \phi)$$

Which gives:

$$f(\theta, \phi) = \sum_{l=0}^{\infty} \sum_{k=-l}^{l} \hat{f}_{lk} \cdot Y^k_l(\theta, \phi)$$
Implications

\[ Y^k_l(\theta, \phi) = e^{ik\theta} \cdot P^{|k|}_l(\cos \phi) \]

When the function \( f \) is real-valued, we may want to express it as the sum of real-valued functions.

We can do this by considering the real and imaginary parts of the harmonics independently:

\[
\text{Re} \left( Y^k_l(\theta, \phi) \right) = \cos(k\theta) \cdot P^{|k|}_l(\cos \phi)
\]

\[
\text{Im} \left( Y^k_l(\theta, \phi) \right) = \sin(k\theta) \cdot P^{|k|}_l(\cos \phi)
\]

Note that since \( Y^k_l = \overline{Y}^{-k}_l \),

\[
\text{Span}\{\text{Re}(Y^k_l), \text{Im}(Y^k_l)\} = \text{Span}\{Y^k_l, Y^{-k}_l\}
\]
The Spherical Harmonics

When the function $f$ is real-valued, we may want to express it as the sum of real-valued functions.

\[ F_0 = \text{Span} \left( \right) \]
\[ F_1 = \text{Span} \left( \right) \]
\[ F_2 = \text{Span} \left( \right) \]
\[ F_3 = \text{Span} \left( \right) \]
Implications

Since the spaces $V^l$ are sub-representations, rotating a function that is the sum of the $l$-th spherical harmonics, will give a function that is the sum of the $l$-th spherical harmonics.
Implications

Since the spaces $V^l$ are sub-representations, rotating a function that is the sum of the $l$-th spherical harmonics, will give a function that is the sum of the $l$-th spherical harmonics.
Implications

Since the spaces $V^l$ are sub-representations, rotating a function that is the sum of the $l$-th spherical harmonics, will give a function that is the sum of the $l$-th spherical harmonics.
Implications

Since the spaces $V^l$ are sub-representations, rotating a function that is the sum of the $l$-th spherical harmonics, will give a function that is the sum of the $l$-th spherical harmonics.
Application (Shape Matching)

Goal:

Given a spherical function representing the surface of a 3D model by a spherical function:

we would like a rotation invariant representation.
Application (Shape Matching)

Approach:

We use the facts that:

- Rotations are unitary.
- The spherical harmonics of degree $l$ are representations of the group of rotation.
Application (Shape Matching)

Approach:

1. Given a spherical function, compute its spherical harmonic decomposition:

\[
f(\theta, \phi) = \sum_{l=0}^{\infty} \sum_{k=-l}^{l} \hat{f}_{lk} \cdot Y_{l}^{k}(\theta, \phi)
\]
Application (Shape Matching)

Approach:

2. Combine the spherical harmonics of the same degree...

\[ f(\theta, \phi) = \sum_{l=0}^{\infty} \sum_{k=-l}^{l} \hat{f}_{lk} \cdot Y_l^k(\theta, \phi) \]

\[ f(\theta, \phi) = \sum_{l=0}^{\infty} f_l(\theta, \phi) \]
Application (Shape Matching)

Approach:

3. Store the norms:
\[ \{ \| f_0 \|, \| f_1 \|, \cdots \} \]
These provide a rotation invariant descriptor.

Norms Invariant to Rotation

\[ f(\theta, \phi) = \sum_{l=0}^{\infty} f_l(\theta, \phi) = \text{Constant} + \text{1st Degree} + \text{2nd Degree} + \text{3rd Degree} + \cdots \]