

FFTs in Graphics and Vision

Correlation and Convolution

Outline



Review

Correlation:

- One-Dimensional (Continuous)
- One-Dimensional (Discrete)
- Higher-Dimensional
- Computational Complexity

Convolution

Applications in 1D

Representations



A <u>representation</u> of a group G <u>on a vector space</u> V, denoted (ρ, V) , is a map ρ that sends every element in G to an invertible linear transformation on V, preserving the group action:

$$\rho(g \cdot h) = \rho(g) \cdot \rho(h) \quad \forall g, h \in G.$$

Sub-Representation



Given a representation (ρ, V) of a group G, if there exists a subspace $W \subset V$ such that the representation maps W to itself (i.e. fixes W): $\rho_g(w) \in W \quad \forall g \in G; \ w \in W$ then we say that W is a sub-representation of V.

Irreducible Representations



Given a representation (ρ, V) of a group G, the representation is said to be <u>irreducible</u> if the only sub-representations of V are:

$$W = V$$
 and $W = \{0\}$

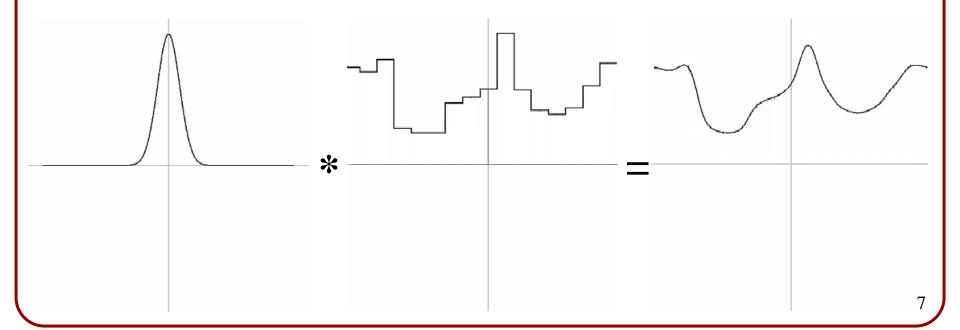
Schur's Lemma (Corollary)



If (ρ, V) is an irreducible, (unitary), representation of a commutative group G, then V must be one-dimensional.

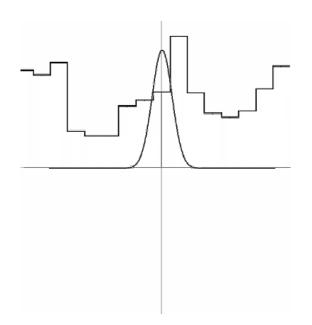


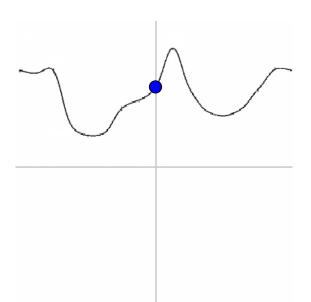
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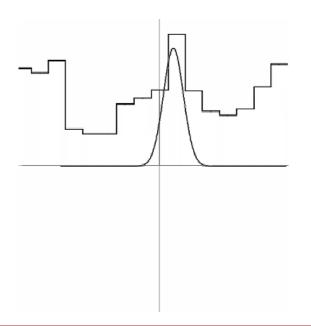
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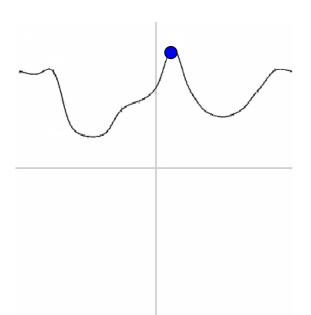






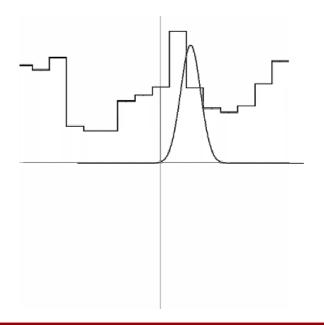
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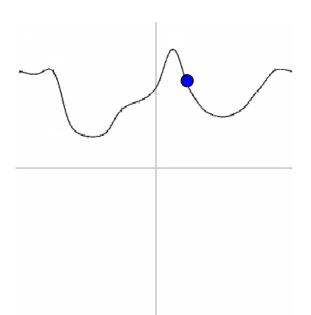






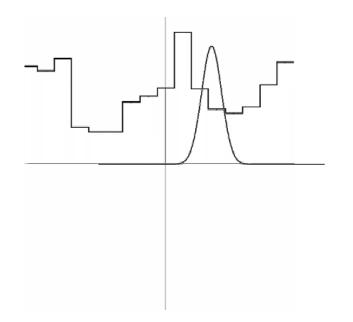
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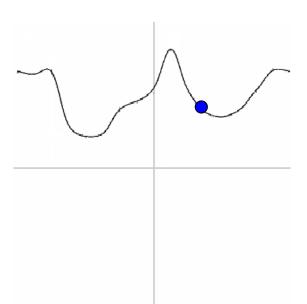






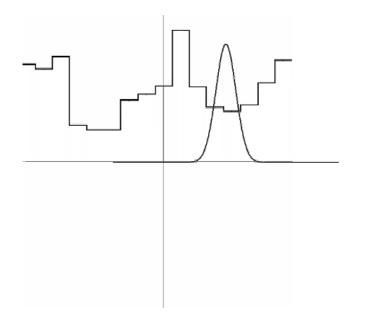
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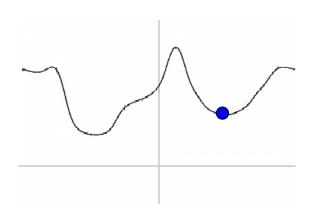






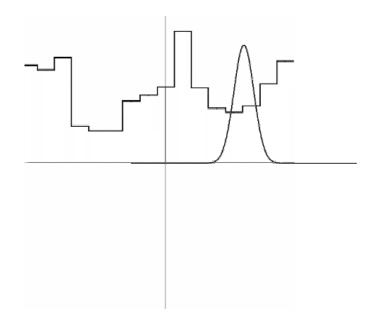
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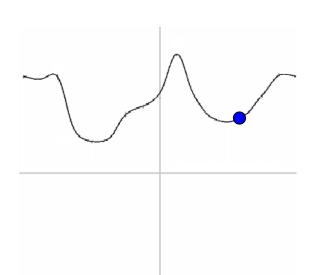






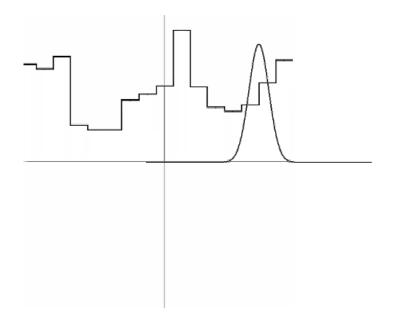
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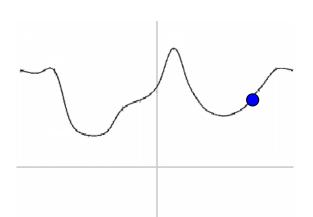




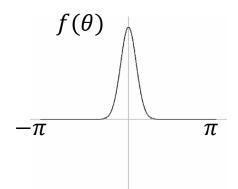


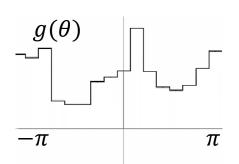
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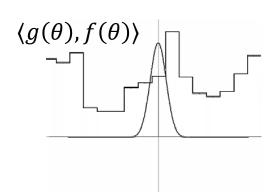




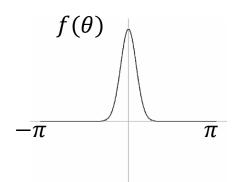


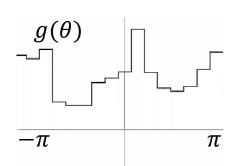


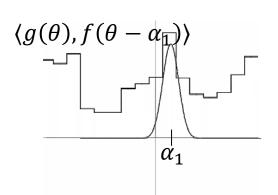




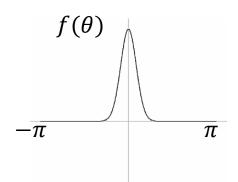


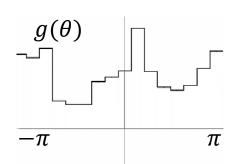


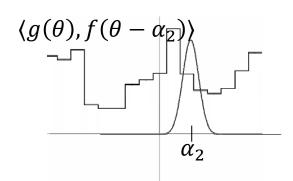




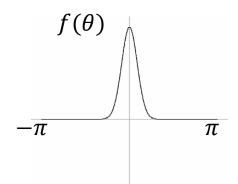


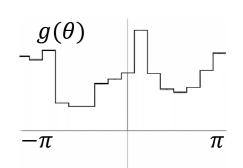


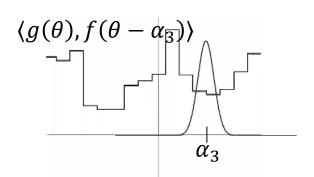










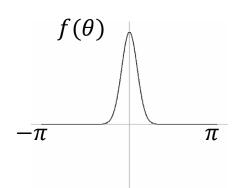


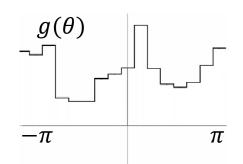


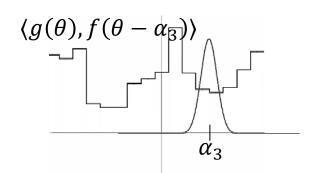
We can write out the operation of smoothing a signal g by a filter f as:

$$(g \star f)(\alpha) = \langle g, \rho_{\alpha}(f) \rangle$$

where ρ_{α} is the linear transformation that translates a periodic function by α .









We can think of this as a representation:

- $V = L^2(\mathbb{R}/2\pi\mathbb{Z})$ is the space of periodic functions on the line
- \circ $G = \mathbb{R}/2\pi\mathbb{Z}$ is the group \mathbb{R} modulo addition by integer multiples of 2π
- \circ ρ_{α} is the representation translating a function by α .

This is a representation of a commutative group...

Warning:

The domain of functions in V and the space G are both parametrized by points in the range $[0,2\pi)$.

• Though the parameters domains are the same, we should think of them as distinct. (The former is the circle S^1 , the latter is the rotation group SO(2).)

$$\zeta^j \in V, \chi^j : G \to \mathbb{C}$$



⇒ There exist orthogonal one-dimensional (complex) subspaces $V_1, \dots, V_n \subset V$ that are the irreducible representations of V.*

Setting $\zeta^j \in V_j$ to be a unit-vector, we know that the group acts on ζ^j by scalar multiplication.

That is, there exist $\chi^j : G \to \mathbb{C}$ s.t.: $\rho_{\alpha}(\zeta^j) = \chi^j(\alpha) \cdot \zeta^j$

Since the ζ^{j} are unit vectors:

$$\chi^{j}(\alpha) = \langle \rho_{\alpha}(\zeta^{j}), \zeta^{j} \rangle$$

*In reality, there are infinitely many such subspaces.

$$\zeta^j \in V, \chi^j : G \to \mathbb{C}$$



 \Rightarrow There exist orthogonal one-dimensional (complex) subspaces $V_1, \dots, V_n \subset V$ that are the irreducible representations of V.*

Setting $\zeta^j \in V_j$ to be a unit-vector, we know that the group acts on ζ^j by scalar multiplication.

That is, there exist $\chi^j: G \to \mathbb{C}$ s.t.:

Note:

Since

Since the V_i are orthogonal, the function basis $\{\zeta^1, \dots, \zeta^n\}$ is orthonormal.

$$\chi'(\alpha) = \langle \rho_{\alpha}(\varsigma'), \varsigma' \rangle$$

^{*}In reality, there are infinitely many such subspaces.

$\zeta^j \in V, \chi^j : G \to \mathbb{C}$



Setting $\zeta^j \in V_j$ to be a unit-vector, we know that the group acts on ζ^j by scalar multiplication:

$$\rho_{\alpha}(\zeta^{j}) = \chi^{j}(\alpha) \cdot \zeta^{j}$$

We can write out vectors $f, g \in V$ in the basis $\{\zeta^1, ..., \zeta^n\}$ as:

$$f = \hat{f}_1 \cdot \boldsymbol{\zeta}^1 + \dots + \hat{f}_n \cdot \boldsymbol{\zeta}^n$$

$$g = \hat{g}_1 \cdot \boldsymbol{\zeta}^1 + \dots + \hat{g}_n \cdot \boldsymbol{\zeta}^n$$

with $\hat{\mathbf{f}}$, $\hat{\mathbf{g}} \in \mathbb{C}^n$.

 $\zeta^j \in V, \chi^j : G \to \mathbb{C}$



Then the correlation can be written as:

$$(g \star f)(\alpha) = \langle g, \rho_{\alpha}(f) \rangle$$

Expanding in the function basis $\{\zeta^1, ..., \zeta^n\}$:

$$(g \star f)(\alpha) = \left| \sum_{j} \hat{g}_{j} \cdot \zeta^{j}, \rho_{\alpha} \left(\sum_{k} \hat{f}_{k} \cdot \zeta^{k} \right) \right|$$

$$\zeta^j \in V, \chi^j : G \to \mathbb{C}$$



$$(g \star f)(\alpha) = \left\langle \sum_{j} \hat{g}_{j} \cdot \zeta^{j}, \rho_{\alpha} \left(\sum_{k} \hat{f}_{k} \cdot \zeta^{k} \right) \right\rangle$$

Using the linearity of ρ_{α} , the (conjugate)-symmetry of the inner-product, and the orthonormality of $\{\zeta^1, \dots, \zeta^n\}$:

$$= \sum_{j,k} \hat{g}_{j} \cdot \bar{f}_{k} \langle \zeta^{j}, \rho_{\alpha}(\zeta^{k}) \rangle$$

$$= \sum_{j,k} \hat{g}_{j} \cdot \bar{f}_{k} \cdot \bar{\chi}^{k}(\alpha) \langle \zeta^{j}, \zeta^{k} \rangle$$

$$= \sum_{i} \hat{g}_{j} \cdot \bar{f}_{j} \cdot \bar{\chi}^{j}(\alpha)$$

 $\zeta^j \in V, \chi^j : G \to \mathbb{C}$



$$(g \star f)(\alpha) = \sum_{j} \hat{g}_{j} \cdot \bar{f}_{j} \cdot \bar{\chi}^{j}(\alpha)$$

This implies that we can compute the correlation by multiplying the coefficients of g by the (conjugate of the) coefficients of f.

Correlation in the spatial domain is multiplication in the frequency domain!

 $\zeta^j \in V, \chi^j : G \to \mathbb{C}$



What is $\chi^{j}(\alpha)$?

Since the representation is unitary, $|\chi^j(\alpha)| = 1$:

$$\exists \tilde{\chi}^j : G \to \mathbb{R}$$
 s.t. $\chi^j(\alpha) = e^{-i\tilde{\chi}^j(\alpha)}$

Since it's a representation:

$$\exists \kappa_j \in \mathbb{R} \quad \text{s.t.} \quad \tilde{\chi}^j(\alpha) = \kappa_j \cdot \alpha$$

And it takes $\alpha = 2\pi$ to the identity:

$$1 = e^{-i\kappa_j 2\pi} \quad \Leftrightarrow \quad k_j \in \mathbb{Z}$$

 $\zeta^j \in V, \chi^j : G \to \mathbb{C}$



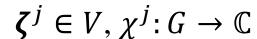
What is $\chi^{j}(\alpha)$?

$$(g \star f)(\alpha) = \sum_{j} \hat{g}_{j} \cdot \bar{f}_{j} \cdot \bar{\chi}^{j}(\alpha)$$

Thus, the correlation of the signals $f, g: S^1 \to \mathbb{C}$ can be expressed as:

$$(g \star f)(\alpha) = \sum_{j} \hat{g}_{j} \cdot \bar{f}_{j} \cdot e^{i\kappa_{j}\alpha}$$

where $\kappa_i \in \mathbb{Z}$.





What is *ζ^j*?

By definition of χ^j , we have:

$$\rho_{\alpha}(\boldsymbol{\zeta}^{j}) = \chi^{j}(\alpha) \cdot \boldsymbol{\zeta}^{j} = e^{-ik_{j}\alpha} \cdot \boldsymbol{\zeta}^{j}$$

for some $k_i \in \mathbb{Z}$.

On the other hand, we have:

$$\boldsymbol{\zeta}^{j}(\theta) \cdot e^{-ik_{j}\alpha} = \left[\rho_{\alpha}(\boldsymbol{\zeta}^{j})\right](\theta) = \boldsymbol{\zeta}^{j}(\theta - \alpha)$$

$$\boldsymbol{\zeta}^{j}(\theta) = c_{j} \cdot e^{ik_{j}\theta}$$

Outline



Review

Correlation:

- One-Dimensional (Continuous)
- One-Dimensional (Discrete)
- Higher-Dimensional
- Computational Complexity

Convolution

Applications in 1D



Let's revisit periodic functions in more detail:

- $V = L^2(\mathbb{R}/2\pi\mathbb{Z})$ is the space of periodic functions on the line
 - $G = \mathbb{R}/2\pi\mathbb{Z}$ is the group \mathbb{R} modulo addition by integer multiples of 2π
- \circ ρ_{α} is the representation translating a function by α :

$$(\rho_{\alpha}(f))(\theta) = f(\theta - \alpha)$$

The irreducible representations are:*

$$V_j = \operatorname{Span}(\tilde{\boldsymbol{\zeta}}^j(\theta) = e^{ik_j\theta}).$$

The corresponding scaling functions are:

$$\chi^{j}(\alpha) = e^{-ik_{j}\alpha}$$
.

^{*}Note that while $\tilde{\zeta}^j$ is an eigenvector of ρ_{α} , it is not unit-norm.



$$V_j = \operatorname{Span}(\tilde{\zeta}^j(\theta) = e^{ik_j\theta})$$

 $\chi^j(\alpha) = e^{-ik_j\alpha}$

The one-dimensional sub-space:

$$\mathrm{Span}\big(\tilde{\boldsymbol{\zeta}}^{j}(\theta) = e^{ik\theta}\big)$$

is a sub-representation for <u>every</u> integer $k \in \mathbb{Z}$.

$$V_j = \operatorname{Span}(\tilde{\zeta}^j(\theta) = e^{ij\theta})$$

 $\chi^j(\alpha) = e^{-ij\alpha}$



Note:

The periodic functions:

$$\tilde{\boldsymbol{\zeta}}^{j}(\theta) = e^{ij\theta}$$

do not have unit norm!

$$\|\tilde{\zeta}^{j}\|^{2} = \int_{0}^{2\pi} e^{ij\theta} \cdot \overline{e^{ij\theta}} d\theta$$
$$= \int_{0}^{2\pi} 1 d\theta$$
$$= 2\pi$$



Note:

The periodic functions:

$$\tilde{\boldsymbol{\zeta}}^{j}(\theta) = e^{ij\theta}$$

do not have unit norm!

⇒ Normalize to make the functions unit-norm:

$$\boldsymbol{\zeta}^{j}(\theta) = \frac{e^{ij\theta}}{\sqrt{2\pi}}$$

$$\chi^{j}(\alpha) = e^{-ij\alpha} = \sqrt{2\pi} \cdot \zeta^{-j}(\alpha)$$



Remark:

The functions $\{\zeta^k\}_{k=-\infty}^{\infty}$ form an orthonormal basis called the Fourier basis.



Given functions, $f, g \in L^2(\mathbb{R}/2\pi\mathbb{Z})$, we can expand them in the basis $\{\zeta^k\}_{k=-\infty}^{\infty}$:

$$f = \sum_{k=-\infty}^{\infty} \hat{f}_k \cdot \zeta^k$$
 and $g = \sum_{k=-\infty}^{\infty} \hat{g}_k \cdot \zeta^k$

This gives:

$$(g \star f)(\alpha) = \sum_{k=-\infty}^{\infty} \hat{g}_k \cdot \bar{\hat{f}}_k \cdot \bar{\chi}^k(\alpha)$$
$$= \sum_{k=-\infty}^{\infty} \sqrt{2\pi} \cdot \hat{g}_k \cdot \bar{\hat{f}}_k \cdot \boldsymbol{\zeta}^k(\alpha)$$

Correlation (Periodic Functions)



$$\rho_{\alpha}(\boldsymbol{\zeta}^{j}) = e^{-ij\alpha} \cdot \boldsymbol{\zeta}^{j}$$

What's really going on here?

If we express a complex number in terms of radius and angle (r, θ) , then rotation by α degrees corresponds to the map:

$$(r,\theta) \to (r,\theta + \alpha)$$

$$\updownarrow$$

$$re^{i\theta} \to re^{i(\theta + \alpha)} = e^{i\alpha} \cdot re^{i\theta}$$

⇒ Multiplication by a complex, unit-norm, number is the same as translation

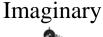
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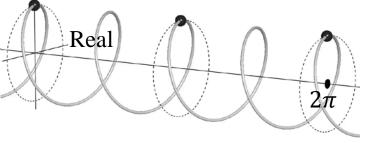


$$\rho_{\alpha}(\zeta^{j}) = e^{-ij\alpha} \cdot \zeta^{j}$$

What's really going on here?

- Visualize complex-valued (periodic) functions the line by drawing the values in the perpendicular plane.
- A complex exponential becomes a helix.
- We can translate the helix along the line.
- This is the same as rotating in the (complex) plane that is perpendicular to the line.





$$f(\theta) = e^{i4\theta}$$

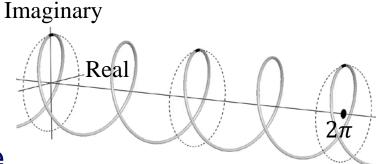
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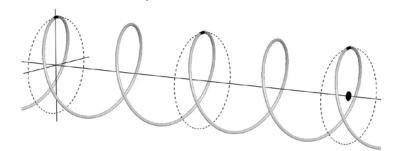
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$$f(\theta) = e^{i4\theta}$$



Translating the domain of a complex exponential



Multiplying its value by a (unit-normal) complex number

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In practice, we don't have infinite precision, and we discretize the function space and the group:

- $V = \mathbb{R}^n$ is the space of periodic n-dimensional arrays
- \circ $G = \mathbb{Z}/n\mathbb{Z}$ is the group of integers modulo n
- \circ ρ_j is the representation shifting the entries in the array by j positions

What are the irreducible representations V_j ?

What are the corresponding scaling functions $\mathbf{x}^j \in \mathbb{C}^n$?

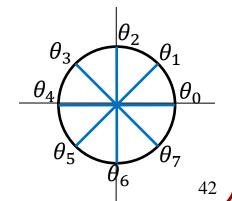


We set V_k to be the (1D) spaces spanned by the <u>discretizations</u> of the complex exponentials:

$$V_k = \operatorname{Span}(\tilde{\mathbf{z}}^k)$$

where $\tilde{\mathbf{z}}^k$ is defined by regularly sampling the k-th complex exponential:

$$\tilde{\mathbf{z}}^k = (e^{ik\theta_0}, \cdots, e^{ik\theta_{n-1}})$$
 with $\theta_j = \frac{2\pi j}{n}$





Applying ρ_i to $\tilde{\mathbf{z}}^k$, we get:

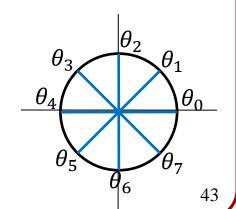
$$\rho_j(\tilde{\mathbf{z}}^k) = \left(e^{ik\theta_{0-j}}, \cdots, e^{ik\theta_{n-1-j}}\right)$$

We can write out:

$$\theta_{m-j} = \frac{2\pi(m-j)}{n}$$

$$= \frac{2\pi m}{n} + \frac{-2\pi j}{n}$$

$$= \theta_m + \theta_{-j}$$





Applying ρ_i to $\tilde{\mathbf{z}}^k$, we get:

$$\rho_j(\tilde{\mathbf{z}}^k) = \left(e^{ik\theta_{0-j}}, \cdots, e^{ik\theta_{n-1-j}}\right)$$

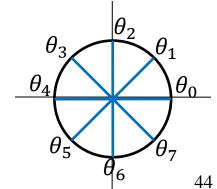
We can write out:

$$\theta_{m-j} = \theta_m + \theta_{-j}$$

Thus:

$$\rho_{j}(\tilde{\mathbf{z}}^{k}) = (e^{ik\theta_{0}} \cdot e^{ik\theta_{-j}}, \cdots, e^{ik\theta_{n-1}} \cdot e^{ik\theta_{-j}})$$
$$= e^{ik\theta_{-j}} \cdot \tilde{\mathbf{z}}^{k}$$

$$\mathbf{x}_i^k = e^{-ik\theta_j}$$





Note 1:

The periodic arrays:

$$\tilde{\mathbf{z}}^k = \left(e^{ik\theta_0}, \cdots, e^{ik\theta_{n-1}}\right)$$

do not have unit norm!

$$\langle \tilde{\mathbf{z}}^k, \tilde{\mathbf{z}}^k \rangle_{[0,2\pi)} = \frac{2\pi}{n} \cdot \sum_{j=0}^{n-1} \tilde{\mathbf{z}}_j^k \cdot \tilde{\tilde{\mathbf{z}}}_j^k$$

$$= \frac{2\pi}{n} \cdot \sum_{j=0}^{n-1} e^{ik\theta_j} \cdot e^{-ik\theta_j} \xrightarrow[\theta_4]{\theta_3} \xrightarrow[\theta_4]{\theta_5}$$

$$= 2\pi$$



Note 1:

The periodic arrays:

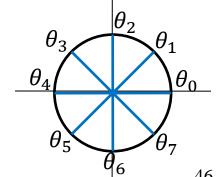
$$\tilde{\mathbf{z}}^k = \left(e^{ik\theta_0}, \cdots, e^{ik\theta_{n-1}}\right)$$

do not have unit norm!

$$\mathbf{x}^k = \sqrt{2\pi} \cdot \mathbf{z}^{-k}$$

We need to normalize these vectors to make them unit-norm:

$$\mathbf{z}^k = \left(\frac{e^{ik\theta_0}}{\sqrt{2\pi}}, \cdots, \frac{e^{ik\theta_{n-1}}}{\sqrt{2\pi}}\right)$$



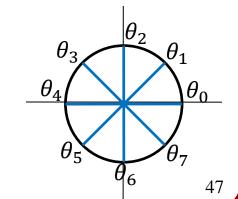


Note 2:

The arrays \mathbf{z}^k and \mathbf{z}^{k+n} are the same array:

$$\mathbf{z}^{k+n} = \left(\frac{e^{i(k+n)\theta_0}}{\sqrt{2\pi}}, \cdots, \frac{e^{i(k+n)\theta_{n-1}}}{\sqrt{2\pi}}\right)$$

$$= \left(\frac{e^{ik\theta_0}}{\sqrt{2\pi}} \cdot e^{in\theta_0}, \cdots, \frac{e^{ik\theta_{n-1}}}{\sqrt{2\pi}} \cdot e^{in\theta_{n-1}}\right)$$





Note 2:

The arrays \mathbf{z}^k and \mathbf{z}^{k+n} are the same array:

$$\mathbf{z}^{k+n} = \left(\frac{e^{ik\theta_0}}{\sqrt{2\pi}} \cdot e^{in\theta_0}, \dots, \frac{e^{ik\theta_{n-1}}}{\sqrt{2\pi}} \cdot e^{in\theta_{n-1}}\right)$$

But $n\theta_i$ is a multiple of 2π :

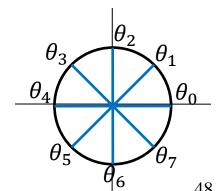
$$n\theta_{j} = \frac{n2\pi j}{n} = 2\pi j$$

$$0$$

$$0$$

$$0$$

$$e^{in\theta_{j}} = 1$$





Note 2:

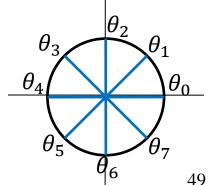
The arrays \mathbf{z}^k and \mathbf{z}^{k+n} are the same array:

$$e^{in\theta_j}=1$$

$$\mathbf{z}^{k+n} = \left(\frac{e^{ik\theta_0}}{\sqrt{2\pi}} \cdot e^{in\theta_0}, \dots, \frac{e^{ik\theta_{n-1}}}{\sqrt{2\pi}} \cdot e^{in\theta_{n-1}}\right)$$

$$= \left(\frac{e^{ik\theta_0}}{\sqrt{2\pi}}, \dots, \frac{e^{ik\theta_{n-1}}}{\sqrt{2\pi}}\right)$$

$$= \mathbf{z}^k$$





Note 3:

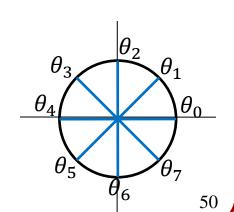
The arrays \mathbf{z}^k and \mathbf{z}^{-k} are conjugate.

$$\mathbf{z}^{-k} = \left(\frac{e^{i(-k)\theta_0}}{\sqrt{2\pi}}, \dots, \frac{e^{i(-k)\theta_{n-1}}}{\sqrt{2\pi}}\right)$$

$$= \left(\frac{e^{-ik\theta_0}}{\sqrt{2\pi}}, \dots, \frac{e^{-ik\theta_{n-1}}}{\sqrt{2\pi}}\right)$$

$$= \left(\frac{e^{ik\theta_0}}{\sqrt{2\pi}}, \dots, \frac{e^{ik\theta_{n-1}}}{\sqrt{2\pi}}\right)$$

$$= \overline{\mathbf{z}}^k$$





Note 4:

Reflecting the coefficients of array z^k through the origin also gives the conjugate:

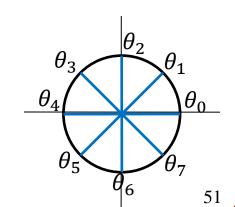
$$z_{-j}^{k} = \frac{e^{ik\theta_{-j}}}{\sqrt{2\pi}}$$

$$= \frac{e^{-ik\theta_{j}}}{\sqrt{2\pi}}$$

$$= \frac{e^{ik\theta_{-j}}}{\sqrt{2\pi}}$$

$$= \frac{e^{ik\theta_{-j}}}{\sqrt{2\pi}}$$

$$= \bar{z}_{j}^{k}$$





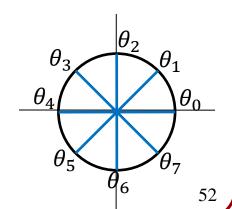
Note 5:

The arrays $\{\mathbf{z}^0, \dots, \mathbf{z}^{n-1}\}$ are linearly independent.

(More specifically, the arrays are orthonormal.)

 \Rightarrow Since the space V is n-dimensional, the arrays form a basis.

This is the <u>discrete Fourier basis</u>.



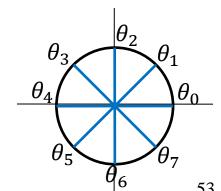


Thus, given two n-dimensional arrays, f and g, we can expand:

$$\mathbf{f} = \sum_{k=0}^{n-1} \hat{f}_k \cdot \mathbf{z}^k \quad \text{and} \quad \mathbf{g} = \sum_{k=0}^{n-1} \hat{g}_k \cdot \mathbf{z}^k$$

This gives:

$$(\mathbf{g} \star \mathbf{f})_{j} = \sum_{k=0}^{n-1} \hat{g}_{k} \cdot \hat{f}_{k} \cdot \bar{x}_{j}^{k}$$
$$= \sqrt{2\pi} \cdot \sum_{k=0}^{n-1} \hat{g}_{k} \cdot \hat{f}_{k} \cdot z_{j}^{k}$$



Outline



Review

Correlation:

- One-Dimensional (Continuous)
- One-Dimensional (Discrete)
- Higher-Dimensional
- Computational Complexity

Convolution

Applications in 1D

Correlation (Higher Dimensions)



The same kind of method can be used for higher dimensions:

Periodic functions in 2D

$$\boldsymbol{\zeta}^{l,m}(\theta,\phi) = \sqrt{\frac{1}{(2\pi)^2}} \cdot e^{il\theta} \cdot e^{im\phi}$$
$$\boldsymbol{\chi}^{l,m}(\alpha,\beta) = \sqrt{(2\pi)^2} \boldsymbol{\zeta}^{-l,-m}(\alpha,\beta)$$

Periodic functions in 3D

$$\boldsymbol{\zeta}^{l,m,n}\left(\theta,\phi,\psi\right) = \sqrt{\frac{1}{(2\pi)^3}} \cdot e^{il\theta} \cdot e^{im\phi} \cdot e^{in\psi}$$
$$\boldsymbol{\chi}^{l,m,n}(\alpha,\beta,\gamma) = \sqrt{(2\pi)^3} \cdot \boldsymbol{\zeta}^{-l,-m,-n}(\alpha,\beta,\gamma)$$

Outline



Review

Correlation:

- One-Dimensional (Continuous)
- One-Dimensional (Discrete)
- Higher-Dimensional
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Convolution

Applications in 1D



To compute the correlation of two periodic, n-dimensional arrays $\mathbf{f}, \mathbf{g} \in \mathbb{C}^n$:

1. Express **f** and **g** in the basis $\{\mathbf{z}^0, ..., \mathbf{z}^{n-1}\} \subset \mathbb{C}^n$:

$$\mathbf{f} = \sum_{k=0}^{n-1} \hat{f}_k \cdot \mathbf{z}^k \quad \text{and} \quad \mathbf{g} = \sum_{k=0}^{n-1} \hat{g}_k \cdot \mathbf{z}^k$$

2. Multiply (and scale) the coefficients:

$$(\mathbf{g} \star \mathbf{f}) = \sqrt{2\pi} \cdot \sum_{k=0}^{n-1} \hat{g}_k \cdot \bar{\hat{f}}_k \cdot \mathbf{z}^k$$

3. Evaluate at every index *j*:

$$(\mathbf{g} \star \mathbf{f})_j = \sqrt{2\pi} \cdot \sum_{k=0}^{n-1} \hat{g}_k \cdot \bar{f}_k \cdot z_j^k$$



To compute the correlation of two periodic, n-dimensional arrays $\mathbf{f}, \mathbf{g} \in \mathbb{C}^n$:

- The first and third steps are a change of bases.
 - Implemented (naively) as matrix multiplication these have complexity $O(n^2)$.



To compute the correlation of two periodic, n-dimensional arrays $\mathbf{f}, \mathbf{g} \in \mathbb{C}^n$:

1. Express **f** and **g** in the basis $\{\mathbf{z}^0, ..., \mathbf{z}^{n-1}\} \subset \mathbb{C}^n$:

$$\mathbf{f} = \sum_{k=0}^{n-1} \hat{f}_k \cdot \mathbf{z}^k \quad \text{and} \quad \mathbf{g} = \sum_{k=0}^{n-1} \hat{g}_k \cdot \mathbf{z}^k$$

 $O(n^2)$

2. Multiply (and scale) the coefficients:

$$(\mathbf{g} \star \mathbf{f}) = \sqrt{2\pi} \cdot \sum_{k=0}^{n-1} \hat{g}_k \cdot \bar{\hat{f}}_k \cdot \mathbf{z}^k$$

O(n)

3. Evaluate at every index *j*:

$$(\mathbf{g} \star \mathbf{f})_j = \sqrt{2\pi} \cdot \sum_{k=0}^{n-1} \hat{g}_k \cdot \bar{f}_k \cdot z_j^k$$

 $O(n^2)$



To compute the correlation of two periodic, n-dimensional arrays $\mathbf{f}, \mathbf{g} \in \mathbb{C}^n$:

The <u>Fast Fourier Transform</u> (FFT) is an algorithm for expressing an array represented by samples at $\{\theta_0, \dots, \theta_{n-1}\}$ as a linear sum of $\{\mathbf{z}^0, \dots, \mathbf{z}^{n-1}\}$.

The <u>Fast Inverse Fourier Transform</u> (IFFT) is an algorithm for expressing an array represented as a linear sum of $\{\mathbf{z}^0, ..., \mathbf{z}^{n-1}\}$ by samples at $\{\theta_0, \dots, \theta_{n-1}\}$.

Both take $O(n \log n)$ time.



To compute the correlation of two periodic, n-dimensional arrays $\mathbf{f}, \mathbf{g} \in \mathbb{C}^n$:

1. Express **f** and **g** in the basis $\{\mathbf{z}^0, ..., \mathbf{z}^{n-1}\} \subset \mathbb{C}^n$:

$$\mathbf{f} = \sum_{k=0}^{n-1} \hat{f}_k \cdot \mathbf{z}^k \quad \text{and} \quad \mathbf{g} = \sum_{k=0}^{n-1} \hat{g}_k \cdot \mathbf{z}^k$$

 $O(n \log n)$

2. Multiply (and scale) the coefficients:

$$(\mathbf{g} \star \mathbf{f}) = \sqrt{2\pi} \cdot \sum_{k=0}^{n-1} \hat{g}_k \cdot \bar{\hat{f}}_k \cdot \mathbf{z}^k$$

O(n)

3. Evaluate at every index *j*:

$$(\mathbf{g} \star \mathbf{f})_j = \sqrt{2\pi} \cdot \sum_{k=0}^{n-1} \hat{g}_k \cdot \bar{f}_k \cdot z_j^k$$

 $O(n \log n)$

Outline



Review

Correlation:

- One-Dimensional (Continuous)
- One-Dimensional (Discrete)
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Applications in 1D



Recall:

The Fourier basis is defined by the functions:

$$\boldsymbol{\zeta}^{j}(\theta) = \frac{e^{ij\theta}}{\sqrt{2\pi}} \quad \forall j \in \mathbb{Z}$$

These functions have the properties that their reflections through the origin are their conjugates:

$$\boldsymbol{\zeta}^{j}(-\theta) = \frac{e^{-ij\theta}}{\frac{\sqrt{2\pi}}{e^{ij\theta}}}$$
$$= \frac{e^{-ij\theta}}{\frac{\sqrt{2\pi}}{e^{ij\theta}}}$$
$$= \frac{\overline{\zeta}^{j}(\theta)}{\overline{\zeta}^{j}(\theta)}$$



Given a function f, the conjugate of its reflection through the origin can be expressed as:

$$\bar{f}(-\theta) = \sum_{j} f_{j} \cdot \zeta^{j}(-\theta)$$

$$= \sum_{j} \bar{f}_{j} \cdot \bar{\zeta}^{j}(-\theta)$$

$$= \sum_{j} \bar{f}_{j} \cdot \bar{\zeta}^{j}(\theta)$$

$$= \sum_{j} \bar{f}_{j} \cdot \zeta^{j}(\theta)$$



Given complex valued functions f and g on the circle, we define the convolution of g with f as the result obtained by

- 1. reflecting and conjugating f, and
- 2. correlating *g* with the transformed *f*:

$$g * f = g * \tilde{f} \quad \text{with } \tilde{f}(\theta) = \bar{f}(-\theta)$$

$$\updownarrow$$

$$g * f = g * \tilde{f} \quad \text{with } \hat{f}_k = \bar{f}_k$$



$$g * f = g * \tilde{f}$$
 with $\hat{f}_k = \bar{f}_k$

Plugging this into equation for correlation:

$$(g * f) = (g * \tilde{f}) = \sqrt{2\pi} \sum_{k} \hat{g}_{k} \cdot \hat{\bar{f}}_{k}^{\bar{k}} \cdot \boldsymbol{\zeta}^{k}$$

$$= \sqrt{2\pi} \sum_{k} \hat{g}_{k} \cdot \hat{\bar{f}}_{k}^{\bar{k}} \cdot \boldsymbol{\zeta}^{k}$$

$$= \sqrt{2\pi} \sum_{k} \hat{g}_{k} \cdot \hat{f}_{k} \cdot \boldsymbol{\zeta}^{k}$$



$$g * f = g * \tilde{f}$$
 with $\hat{f}_k = \bar{f}_k$

Note:

1. Unlike correlation, convolution is symmetric:

$$g * f = f * g$$

2. If *f* is real and symmetric with respect to reflection, then convolution and correlation are the same thing:

$$g * f = g * f$$

$$= \sqrt{2\pi} \sum_{k} \hat{g}_k \cdot \hat{f}_k \cdot \boldsymbol{\zeta}^k$$

Outline



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Correlation:

- One-Dimensional (Continuous)
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Correlation

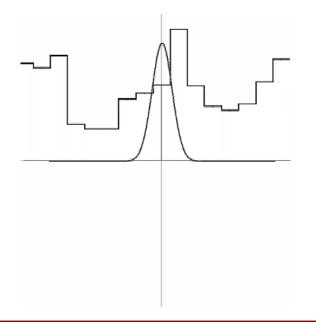
$$f_{-2} \longleftrightarrow g_{-2}$$

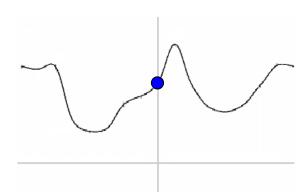
$$f_{-1} \longleftrightarrow g_{-1}$$

$$f_{0} \longleftrightarrow g_{0}$$

$$f_{1} \longleftrightarrow g_{1}$$

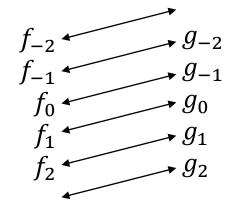
$$f_{2} \longleftrightarrow g_{2}$$

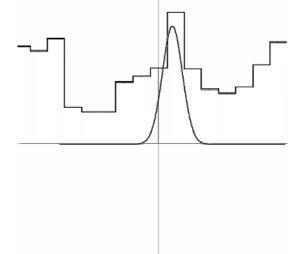


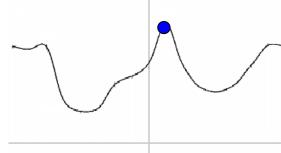




Correlation

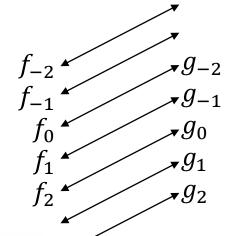


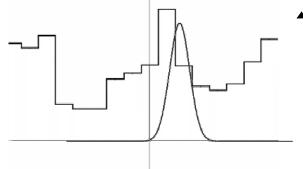


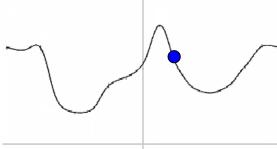




Correlation







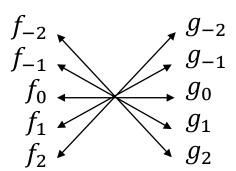


- Correlation
- Convolution

This is like correlation, except that we flip the entries of the array **f** before correlating.

Correlation

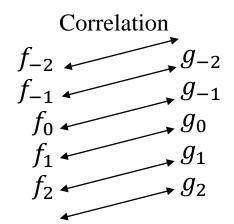
Convolution



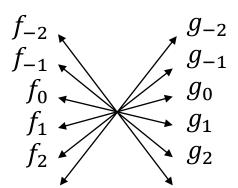


- Correlation
- Convolution

This is like correlation, except that we flip the entries of the array f before correlating.



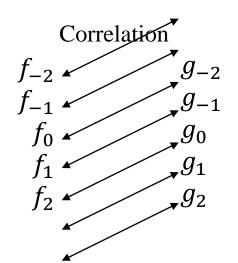
Convolution



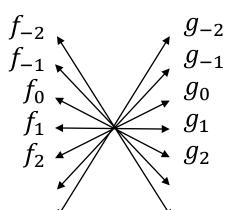


- Correlation
- Convolution

This is like correlation, except that we flip the entries of the array f before correlating.



Convolution





- Correlation
- Convolution
- Polynomial Multiplication



- Correlation
- Convolution
- Polynomial Multiplication

Given two polynomials:

$$p(x) = a_0 + a_1 x + \dots + a_n x^n$$

$$q(x) = b_0 + b_1 x + \dots + b_n x^n$$

we can represent the polynomials p(x) and q(x) by (2n+1)-dimensional arrays:

$$p(x) \to (a_{-n}, \dots, a_{-1}, a_0, a_1, \dots, a_n)$$

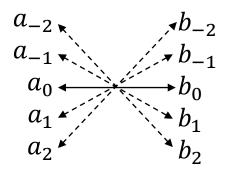
 $q(x) \to (b_{-n}, \dots, b_{-1}, b_0, b_1, \dots, b_n)$

with:

$$a_{-n} = \cdots = a_{-1} = b_{-n} = \cdots = b_{-1} = 0$$

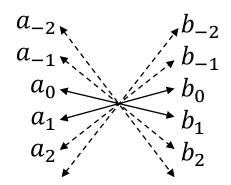


- Correlation
- Convolution
- Polynomial Multiplication
 The 0th order coefficient of the product is:



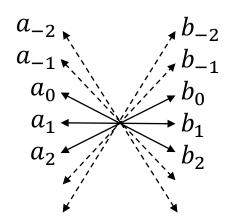


- Correlation
- Convolution
- Polynomial Multiplication
 The 1st order coefficient of the product is:





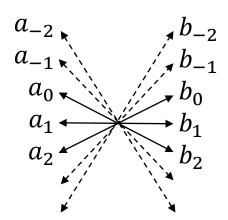
- Correlation
- Convolution
- Polynomial Multiplication
 The 2nd order coefficient of the product is:





- Correlation
- Convolution
- Polynomial Multiplication

The coefficients of the product can be computed efficiently by convolving the arrays corresponding to the coefficients of the original polynomials.





- Correlation
- Convolution
- Polynomial Multiplication
- Big Integer Multiplication
 Given an integer, we can treat it as a polynomial.

$$47601345 = 5 \cdot 10^{0} + 4 \cdot 10^{1} + 3 \cdot 10^{2} + 1 \cdot 10^{3} + \cdots$$



- Correlation
- Convolution
- Polynomial Multiplication
- Big Integer Multiplication

Given an integer, we can treat it as a polynomial. To multiply two integers, we need to figure out what the new value in the 1s place,

$$47601345 = 5 \cdot 10^{0} + 4 \cdot 10^{1} + 3 \cdot 10^{2} + 1 \cdot 10^{3} + \cdots$$

$$46018729 = 9 \cdot 10^{0} + 2 \cdot 10^{1} + 7 \cdot 10^{2} + 8 \cdot 10^{3} + \cdots$$



- Correlation
- Convolution
- Polynomial Multiplication
- Big Integer Multiplication

Given an integer, we can treat it as a polynomial. To multiply two integers, we need to figure out what the new value in the 1s place, the 10s place,

$$47601345 = 5 \cdot 10^{0} + 4 \cdot 10^{1} + 3 \cdot 10^{2} + 1 \cdot 10^{3} + \cdots$$

$$46018729 = 9 \cdot 10^{0} + 2 \cdot 10^{1} + 7 \cdot 10^{2} + 8 \cdot 10^{3} + \cdots$$



- Correlation
- Convolution
- Polynomial Multiplication
- Big Integer Multiplication

Given an integer, we can treat it as a polynomial. To multiply two integers, we need to figure out what the new value in the 1s place, the 10s place, the 10s place, etc. will be.

$$47601345 = 5 \cdot 10^{0} + 4 \cdot 10^{1} + 3 \cdot 10^{2} + 1 \cdot 10^{3} + \cdots$$

$$46018729 = 9 \cdot 10^{0} + 2 \cdot 10^{1} + 7 \cdot 10^{2} + 8 \cdot 10^{3} + \cdots$$



- Correlation
- Convolution
- Polynomial Multiplication
- Big Integer Multiplication

Given an integer, we can treat it as a polynomial. To multiply two integers, we need to figure out what the new value in the 1s place, the 10s place, the 10s place, etc. will be.

So big integer multiplication can be implemented efficiently as a convolution.