



# FFTs in Graphics and Vision

Moving Dot Products



# Outline

## Review

## Moving Dot Products:

- One-Dimensional (Continuous)
- One-Dimensional (Discrete)
- Higher-Dimensional
- Computational Complexity



# Representations

A representation of a group  $G$  on a vector space  $V$  is a map  $\rho$  that sends every element in  $G$  to an invertible linear transformation on  $V$ , satisfying:

$$\rho(g \cdot h) = \rho(g) \cdot \rho(h)$$

for all  $g, h \in G$ .



# Sub-Representation

Given a representation,  $\rho$ , of a group,  $G$ , on a vector space,  $V$ , if there exists a subspace  $W \subset V$ , such that the representation fixes  $W$ :

$$\rho_g w \in W \quad \forall g \in G \text{ and } w \in W$$

then we say that  $W$  is a sub-representation of  $V$ .



# Irreducible Representations

Given a representation,  $\rho$ , of a group,  $G$ , on a vector space,  $V$ , the representation is said to be irreducible if the only subspaces of  $V$  that are sub-representations are:

$$W = V \quad \text{and} \quad W = \emptyset$$



# Schur's Lemma (Corollary)

If  $\rho$  is an irreducible, unitary, representation of a commutative group  $G$  onto a complex vector space  $V$ , then:

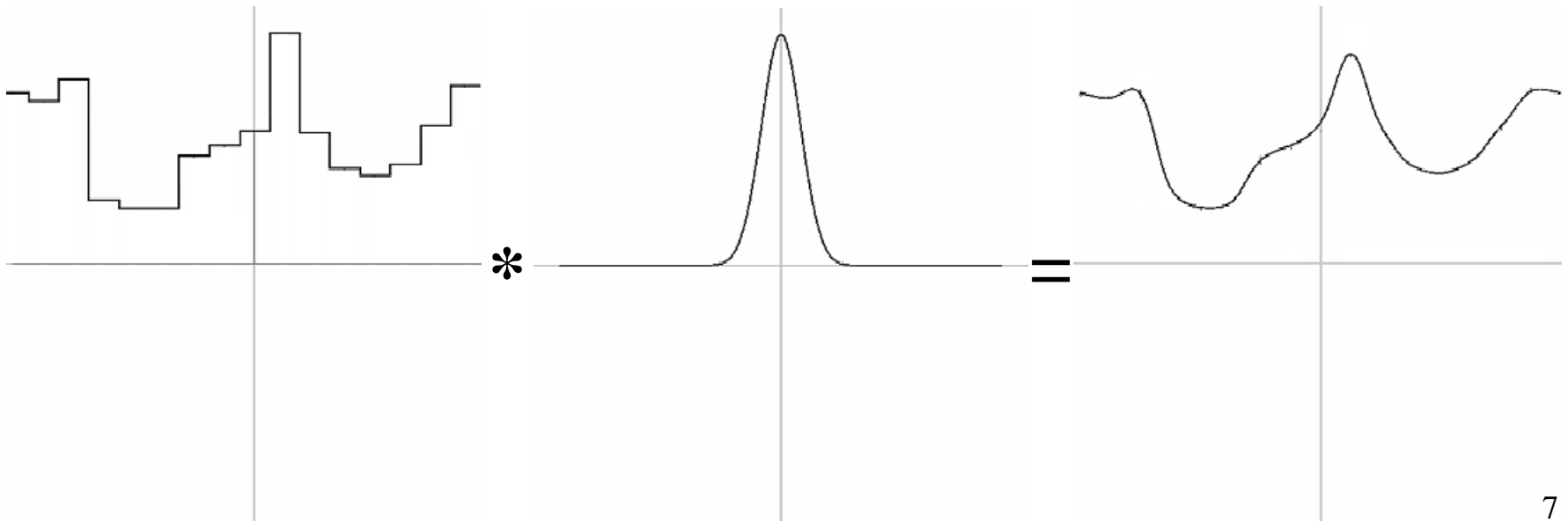
- $V$  must be one-dimensional
- For any  $g \in G$ ,  $\rho(g)$  is a unit-norm complex number



# Why do we care?

In signal/image/voxel processing, we are often interested in applying a filter to some initial data.

E.g. Smoothing:

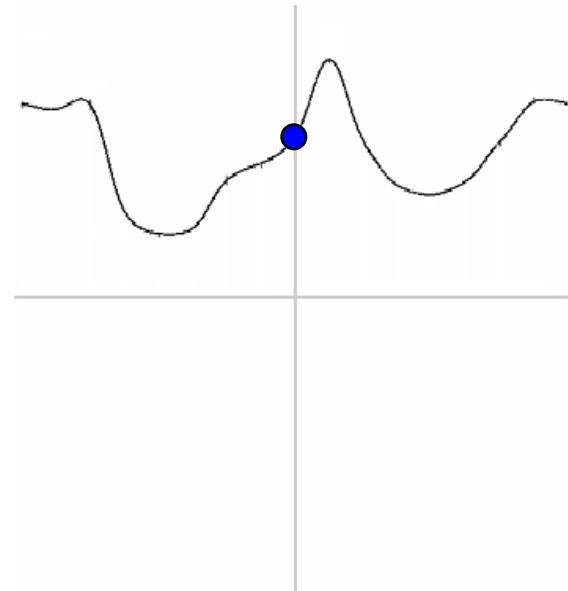
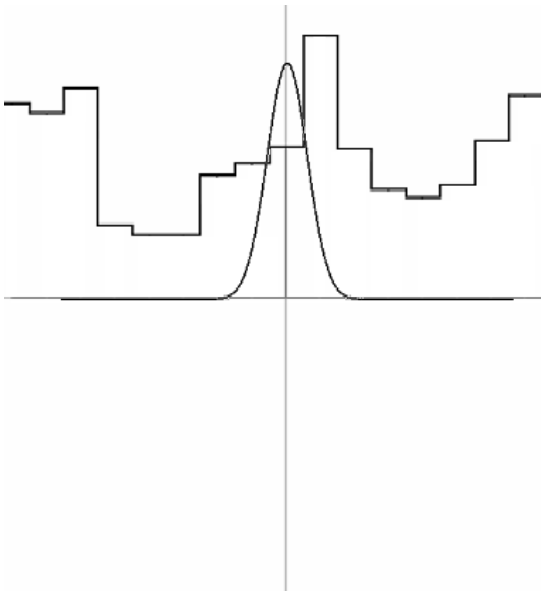




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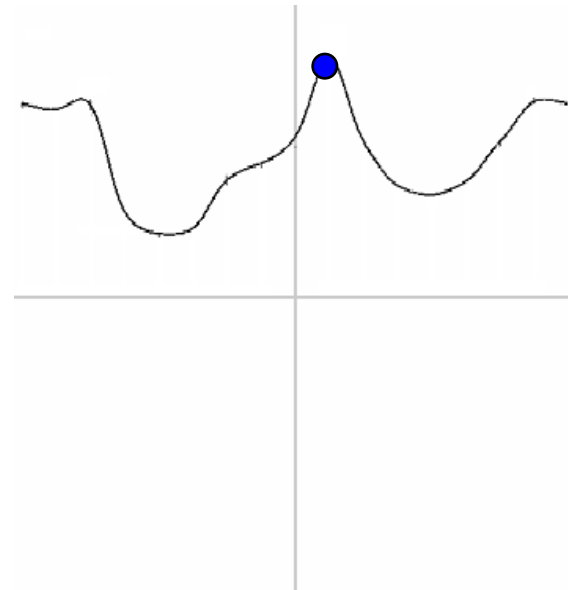
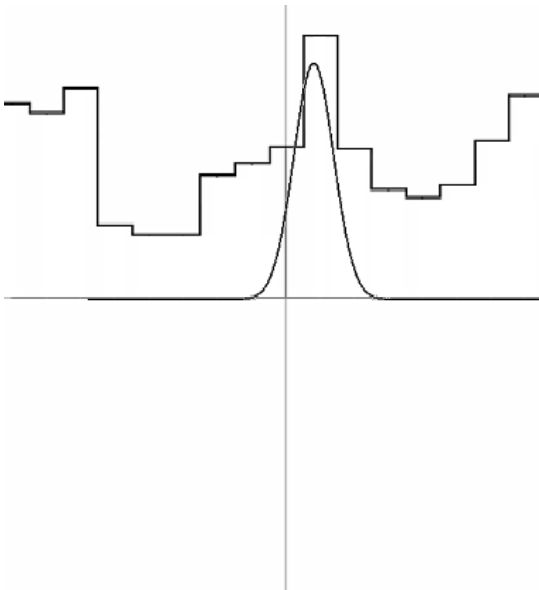




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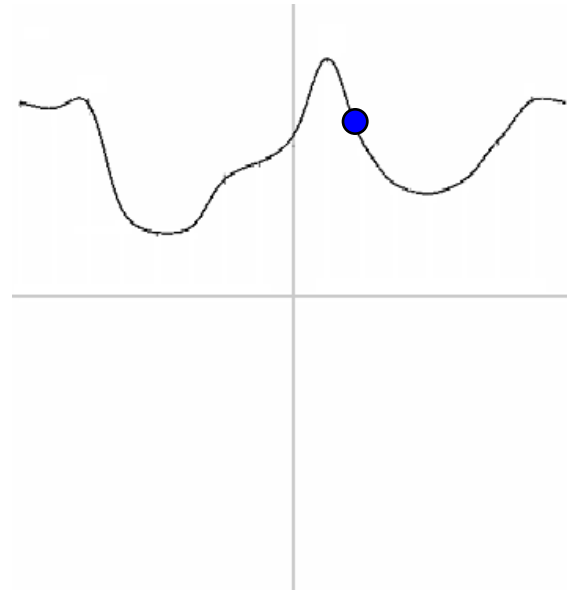
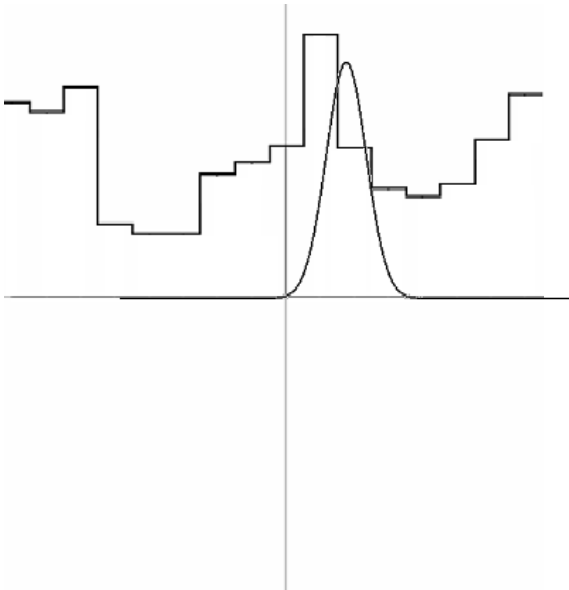




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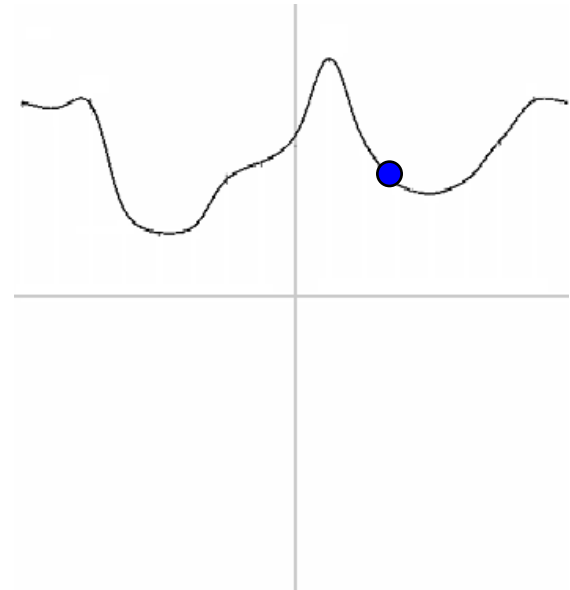
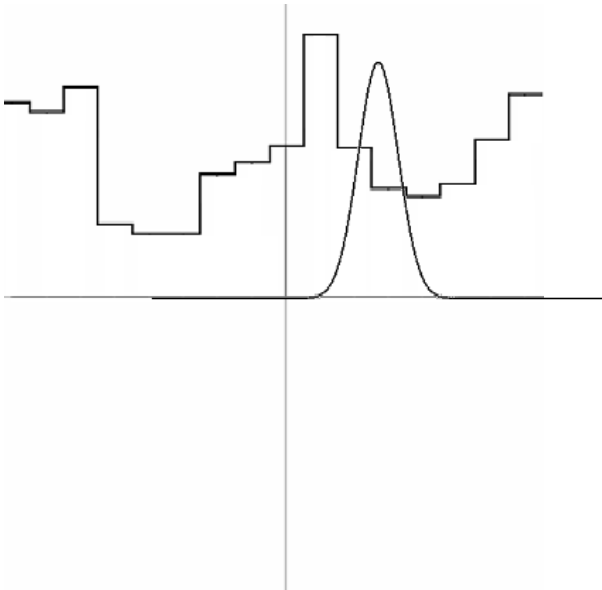




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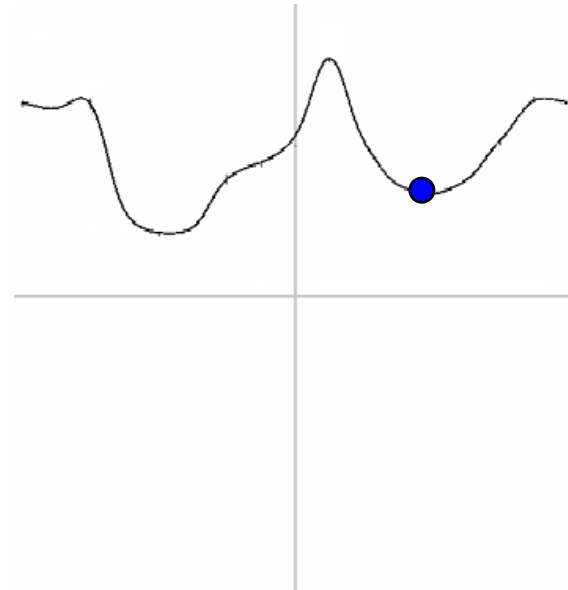
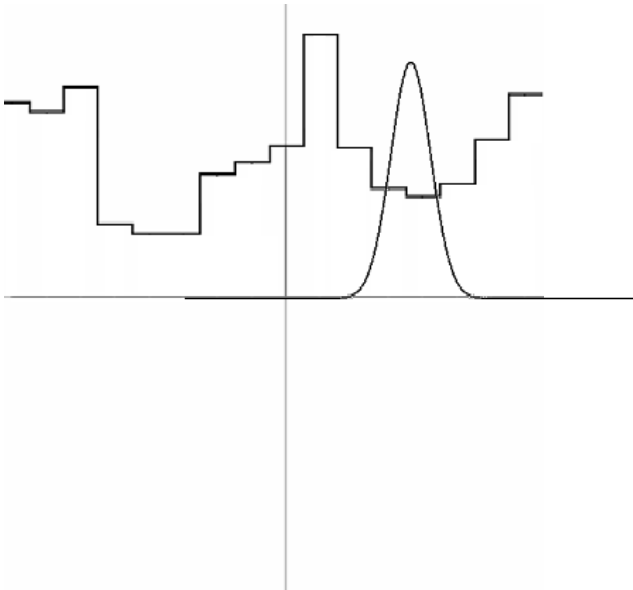




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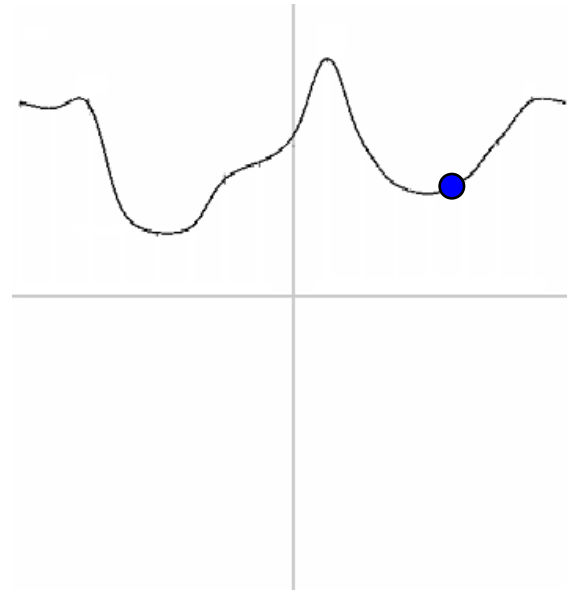
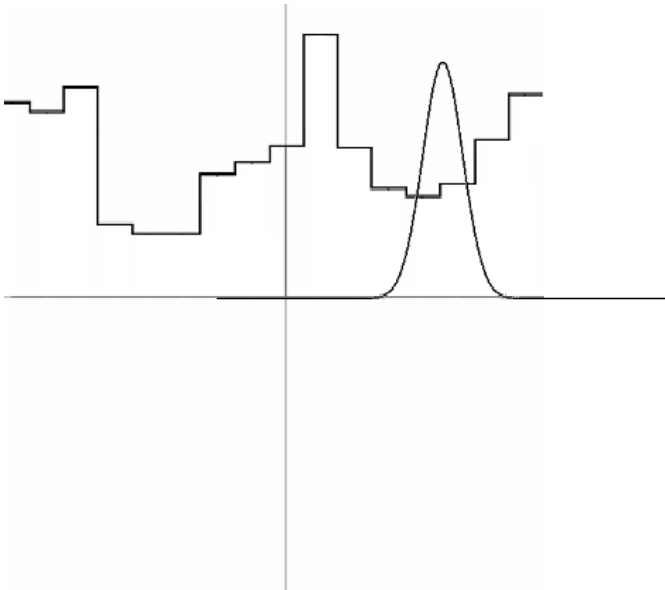




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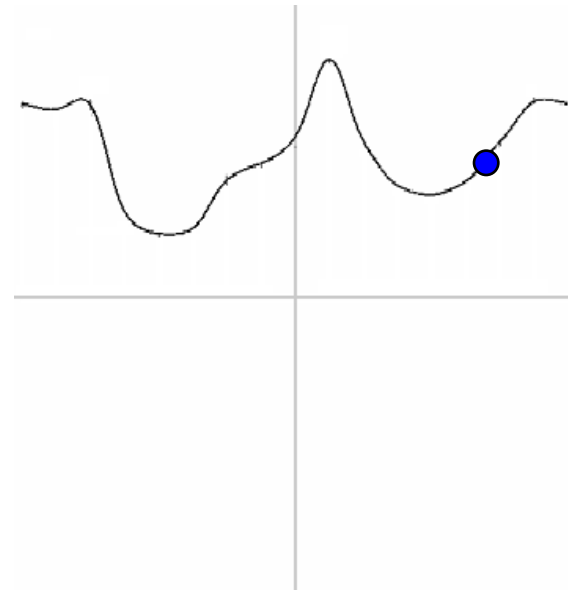
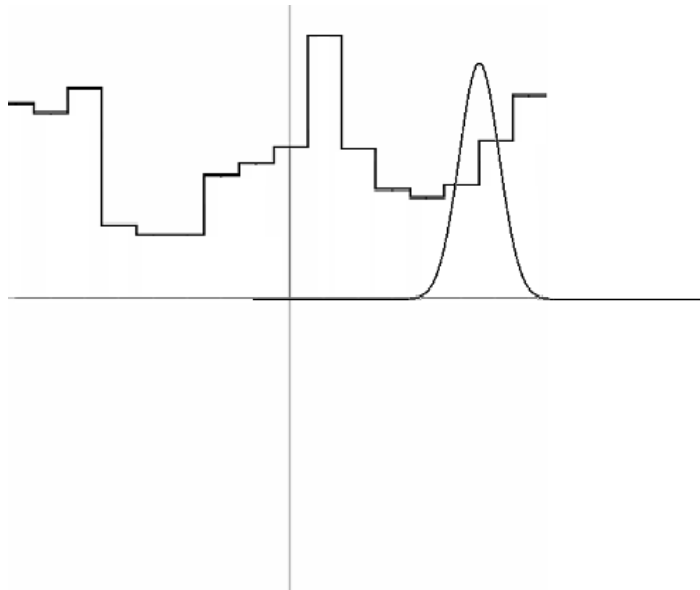




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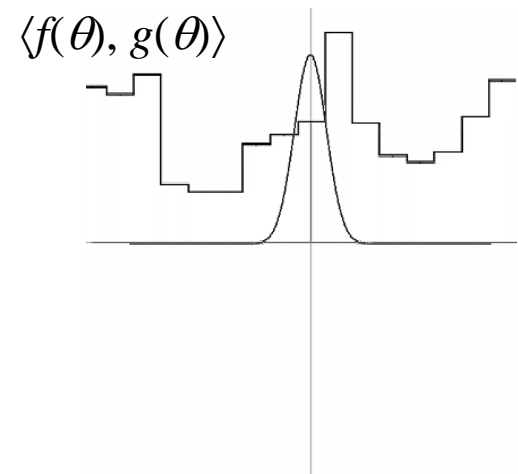
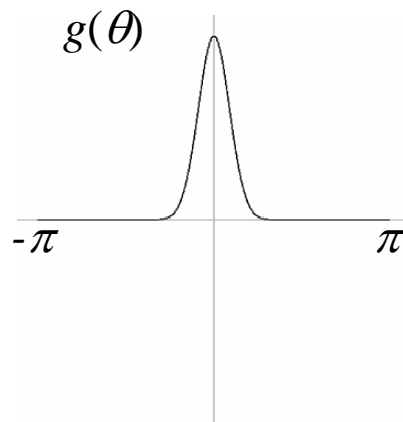
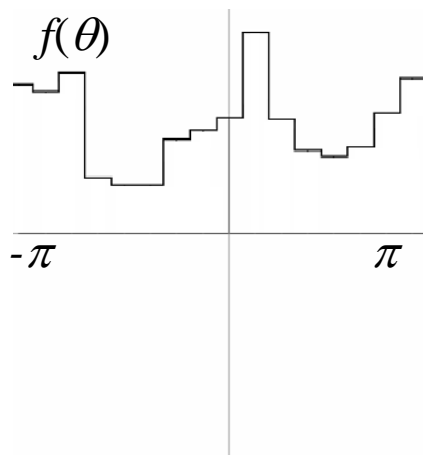
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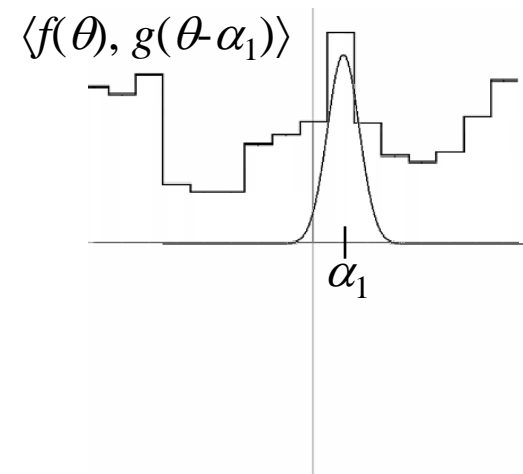
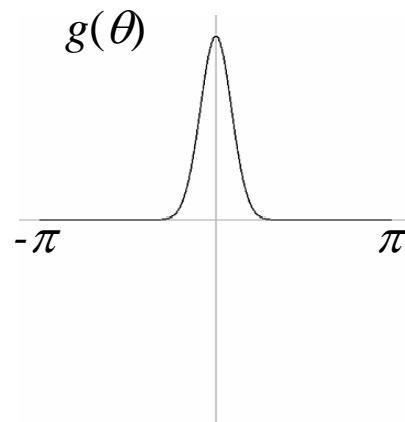
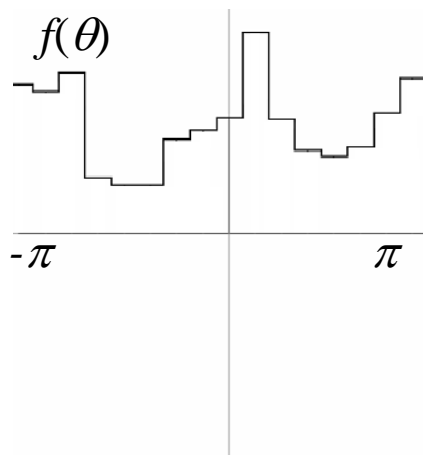
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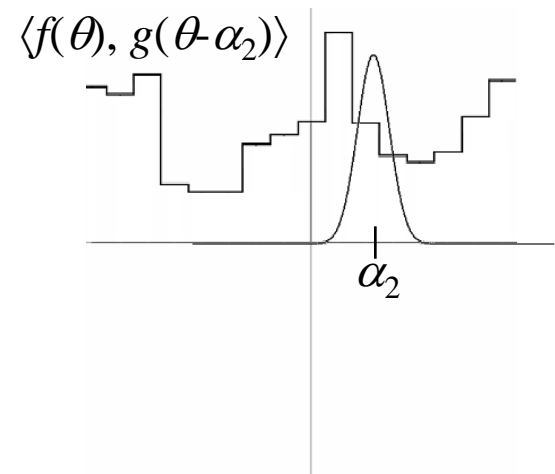
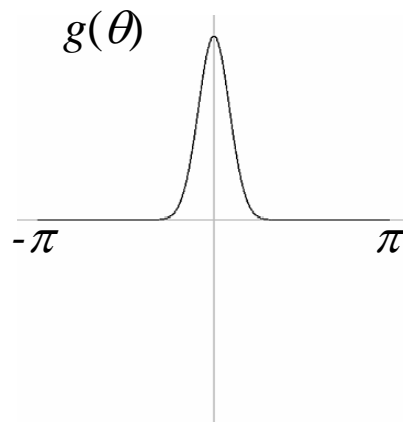
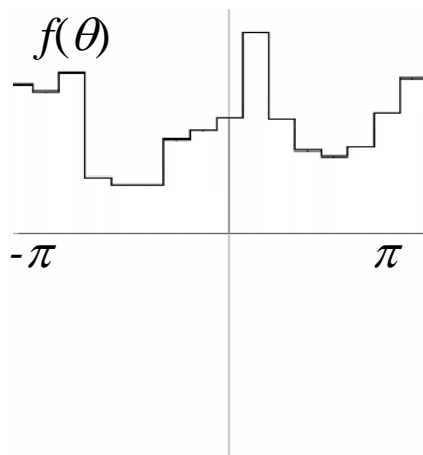
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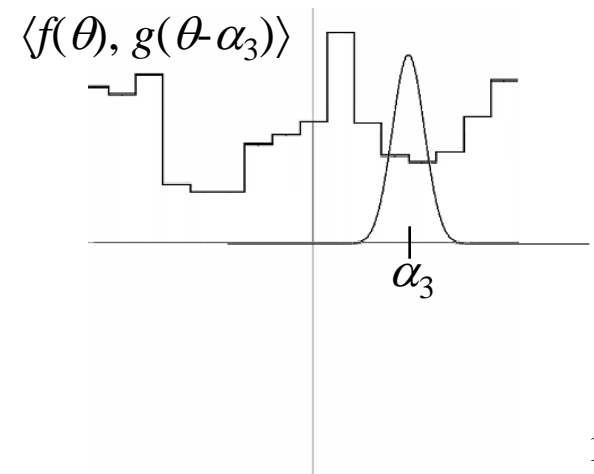
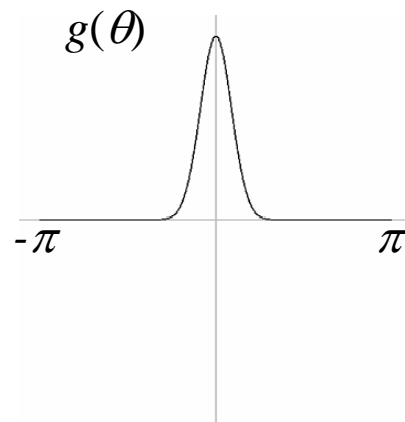
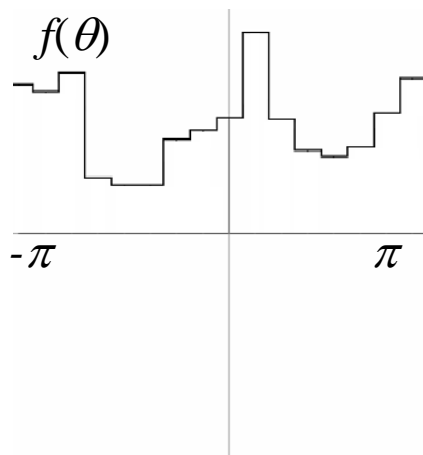
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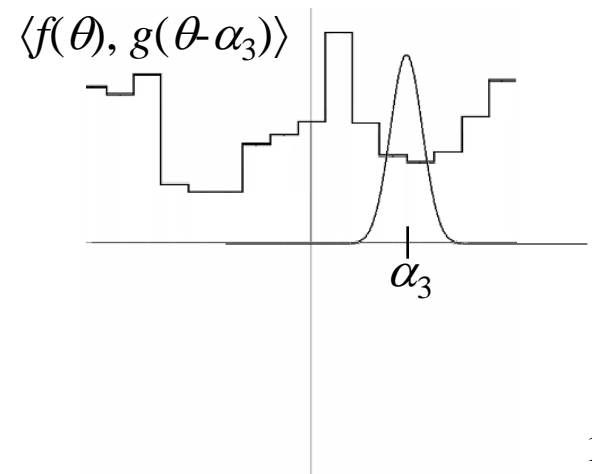
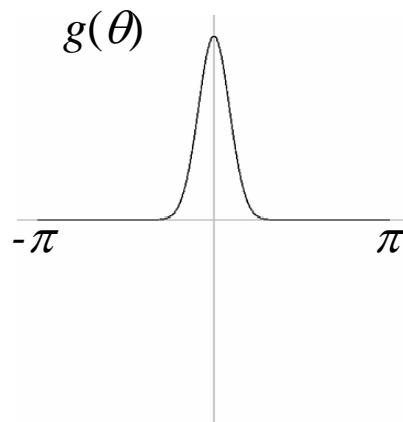
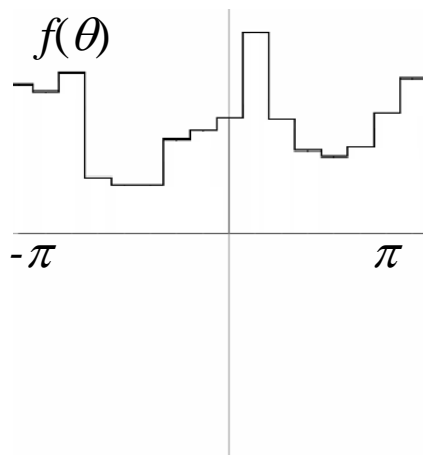


# Smoothing

We can write out the operation of smoothing a signal  $f$  by a filter  $g$  as:

$$(f * g)(\alpha) = \langle f, \rho_\alpha(g) \rangle$$

where  $\rho_\alpha$  is the linear transformation that translates a periodic function by  $\alpha$ .





# Moving Dot Products

We can think of this as a representation:

- $V$  is the space of periodic functions on the line
- $G$  is the group of real numbers in  $[-\pi, \pi)$
- $\rho_\alpha$  is the representation translating a function by  $\alpha$ .



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This is a representation of a commutative group...



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Since the  $V_i$  are sub-representations and they are one-dimensional, we know that:

$$\rho_\alpha(f_i(\theta)) = \lambda_i(\alpha) \cdot f_i(\theta)$$



# Moving Dot Products

Since the  $\{f_i(\theta)\}$  are a basis for  $V$  we can express the functions  $f(\theta)$  and  $g(\theta)$  in terms of this basis:

$$f(\theta) = a_1 f_1(\theta) + a_2 f_2(\theta) + \cdots + a_n f_n(\theta)$$

$$g(\theta) = b_1 f_1(\theta) + b_2 f_2(\theta) + \cdots + b_n f_n(\theta)$$



# Moving Dot Products

Then the moving dot-product can be written as:

$$(f * g)(\alpha) = \langle f, \rho_{\alpha}(g) \rangle$$



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Expanding  $f$  and  $g$  in terms of the basis  $\{f_1, \dots, f_n\}$ :

$$(f * g)(\alpha) = \left\langle \sum_{i=1}^n a_i f_i, \rho_\alpha \left( \sum_{j=1}^n b_j f_j \right) \right\rangle$$



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Using the fact that  $\rho_{\alpha}$  is a linear transformation:

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Using the fact that the inner product is conjugate-linear in the second term:

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Using the fact that on  $V_j$ , the representation  $\rho_\alpha$  is just scalar multiplication:

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And finally, using the fact that the  $f_i$  are orthogonal unit-vectors:

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Convolution in the spatial domain is multiplication in the frequency domain!



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## Moving Dot Products:

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# Moving Dot Products (Periodic Functions)



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- $V$  is the space of periodic functions on the line
- $G$  is the group of real numbers in  $[0, 2\pi)$
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What are the irreducible representations  $V_k$ ?

What are the corresponding functions  $\lambda_k(\alpha)$ ?

# Moving Dot Products (Periodic Functions)



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$$\bar{\lambda}_k(\alpha) = e^{ik\alpha}$$

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We need to normalize these functions to make them unit-norm:

$$f_k(\theta) = \sqrt{\frac{1}{2\pi}} e^{ik\theta}$$



# Moving Dot Products (Periodic Functions)

Thus, given two periodic functions on the line,  $f(\theta)$  and  $g(\theta)$ , we can write:

$$f(\theta) = \sum_{k=-\infty}^{\infty} a_k \sqrt{\frac{1}{2\pi}} e^{ik\theta} \quad \text{and} \quad g(\theta) = \sum_{k=-\infty}^{\infty} b_k \sqrt{\frac{1}{2\pi}} e^{ik\theta}$$

to get:

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If we express a complex number in terms of radius and angle  $(r, \theta)$ , then rotation by  $\alpha$  degrees corresponds to the map:

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$$re^{i\theta} \rightarrow re^{i(\theta+\alpha)} = e^{i\alpha} (re^{i\theta})$$



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What is a rotation by  $\alpha$  degrees in the complex plane?

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Rotating in the complex plane is the same thing as multiplying by a complex, unit-norm, number.

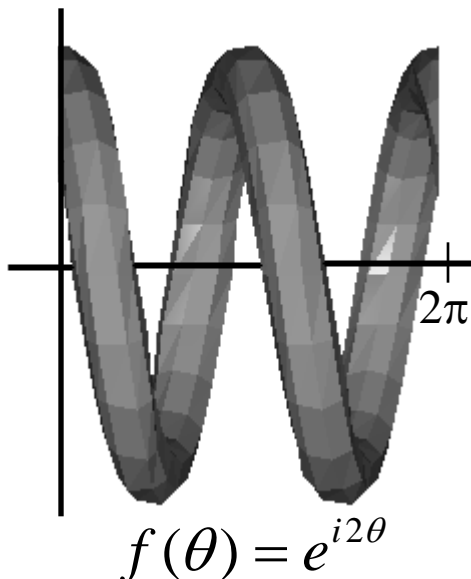
$$re^{i\theta} \rightarrow re^{i(\theta+\alpha)} = e^{i\alpha} (re^{i\theta})$$



# Moving Dot Products (Periodic Functions)

What's really going on here?

Let's consider the graph of a complex exponential. This is just a helix:



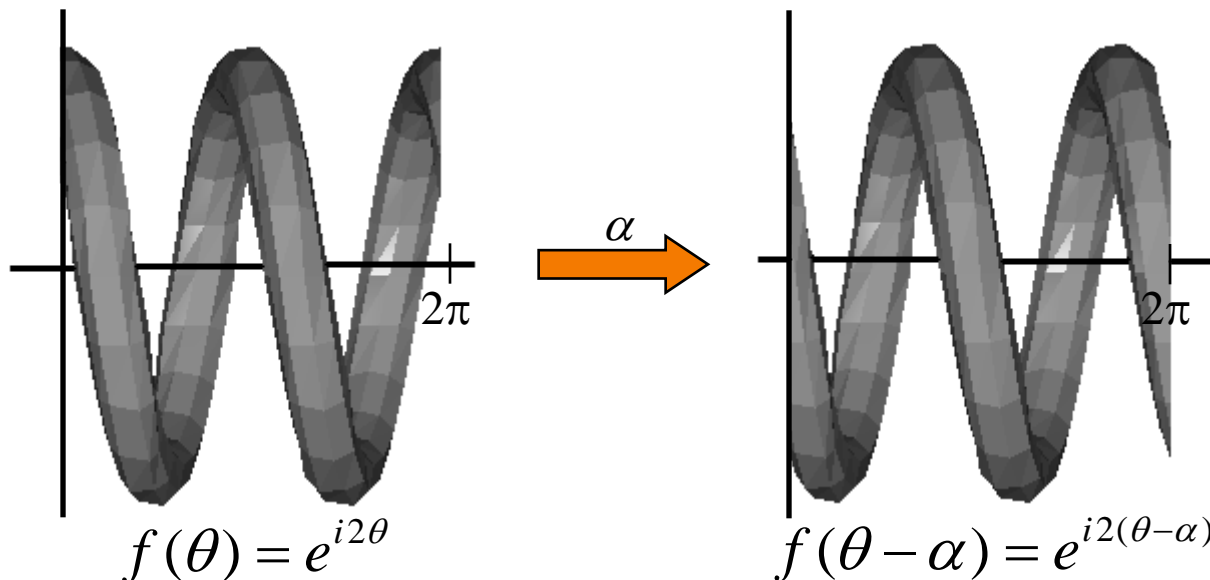


# Moving Dot Products (Periodic Functions)

What's really going on here?

Let's consider the graph of a complex exponential. This is just a helix.

If we translate the function by  $\alpha$ , we get:



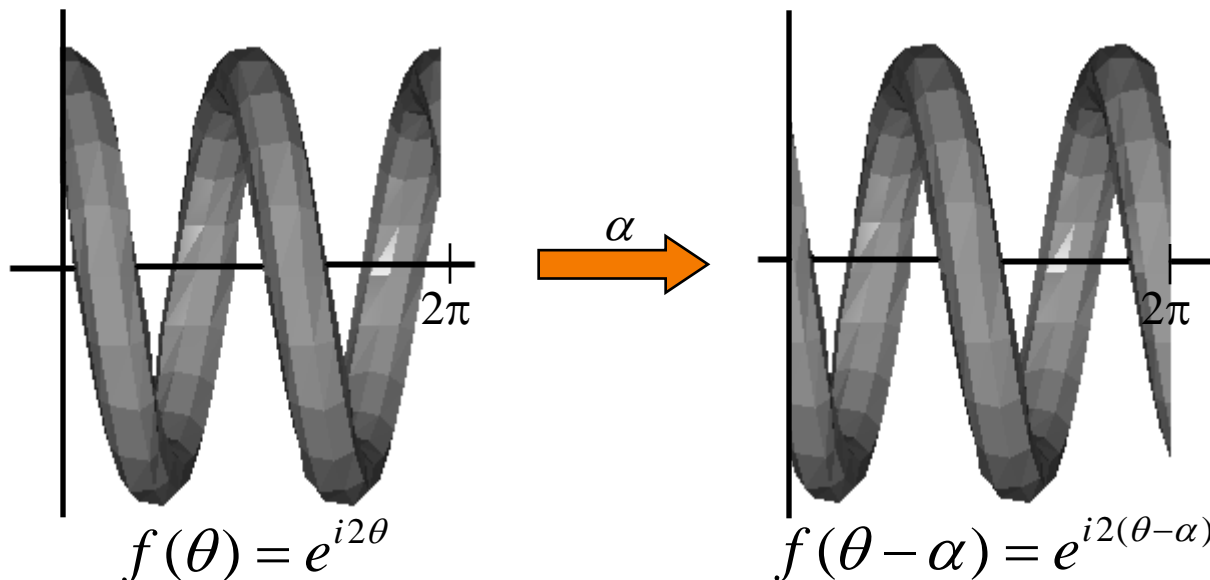


# Moving Dot Products (Periodic Functions)

What's really going on here?

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If we translate a periodic helix is the same thing as rotating it.





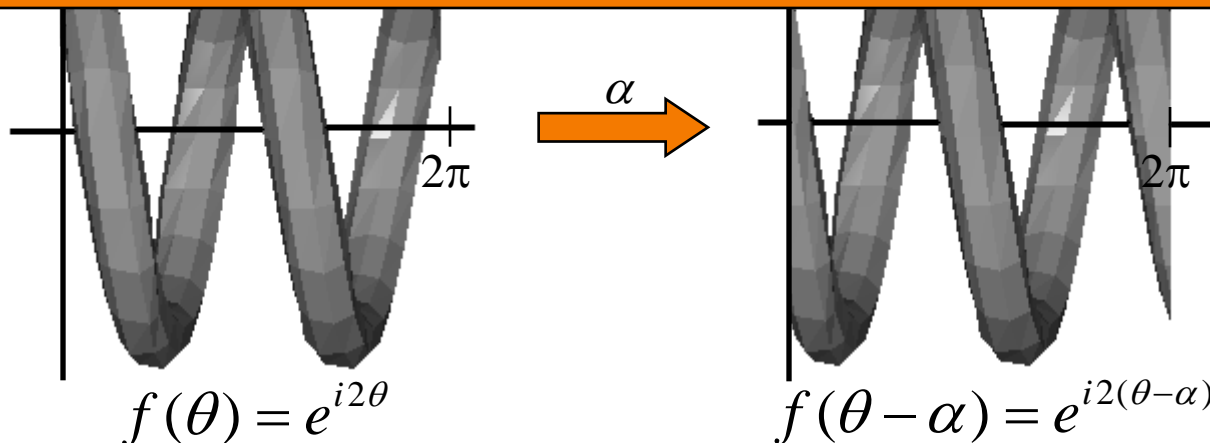
# Moving Dot Products (Periodic Functions)

What's really going on here?

Let's consider the graph of a complex exponential. This is just a helix.

If we translate a periodic helix is the same thing as rotating it.

Rotation in the complex plane is the same thing as multiplication by a complex, unit-norm, number.





# Outline

## Review

## Moving Dot Products:

- One-Dimensional (Continuous)
- **One-Dimensional (Discrete)**
- Higher-Dimensional
- Computational Complexity

# Moving Dot Products (Periodic Arrays)



In practice, we don't have infinite precision, and so we discretize both the function space and the group:

- $V$  is the space of periodic  $n$ -dimensional arrays
- $G$  is the group of integers modulo  $n$
- $\rho_j$  is the representation shifting the entries in the array by  $j$  positions

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- $\rho_j$  is the representation shifting the entries in the array by  $j$  positions

What are the irreducible representations  $V_k$ ?

What are the corresponding functions  $\lambda_k(\alpha)$ ?



# Moving Dot Products (Periodic Arrays)

We set  $V_k$  to be the one-dimensional spaces that are the discretized versions of the complex exponentials:

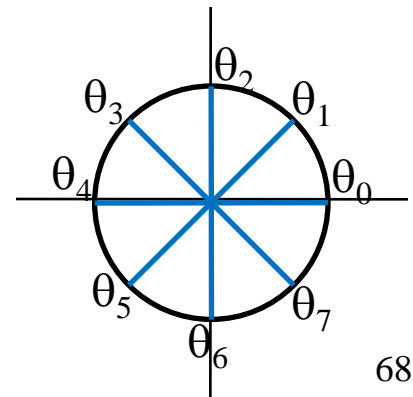
$$V_k = \text{Span}(v_k)$$

where  $v_k$  is defined by sampling the  $k$ -th complex exponential:

$$v_k[\ ] = (e^{ik\theta_0}, e^{ik\theta_1}, \dots, e^{ik\theta_{n-2}}, e^{ik\theta_{n-1}})$$

where:

$$\theta_j = \frac{j2\pi}{n}$$

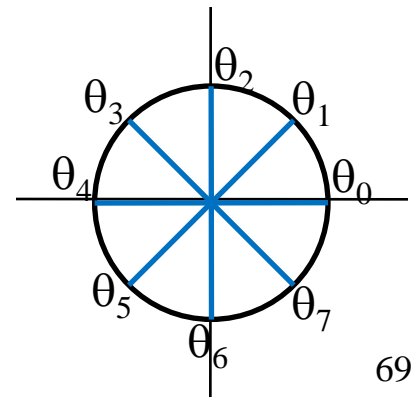


# Moving Dot Products (Periodic Arrays)



Applying  $\rho_\alpha$  to  $v_k[ ]$ , we get:

$$\rho_\alpha(v_k[ ]) = (e^{ik\theta_0-\alpha}, \dots, e^{ik\theta_{n-1}-\alpha})$$



# Moving Dot Products (Periodic Arrays)

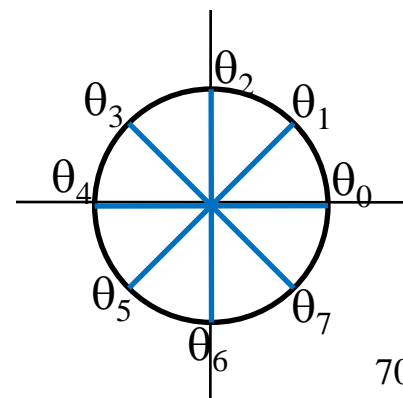


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$$\rho_\alpha(v_k[ ]) = (e^{ik\theta_{0-\alpha}}, \dots, e^{ik\theta_{n-1-\alpha}})$$

Now we can write out:

$$\theta_{j-\alpha} = \frac{(j-\alpha)2\pi}{n}$$



# Moving Dot Products (Periodic Arrays)

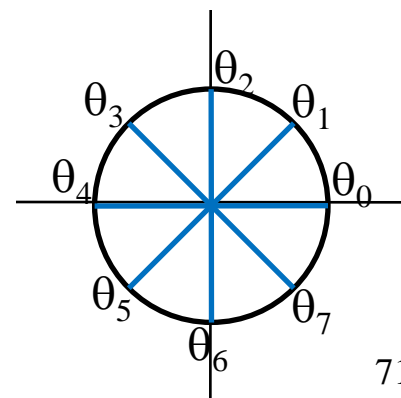


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$$\begin{aligned}\theta_{j-\alpha} &= \frac{(j-\alpha)2\pi}{n} \\ &= \frac{j2\pi}{n} + \frac{-\alpha 2\pi}{n}\end{aligned}$$



# Moving Dot Products (Periodic Arrays)

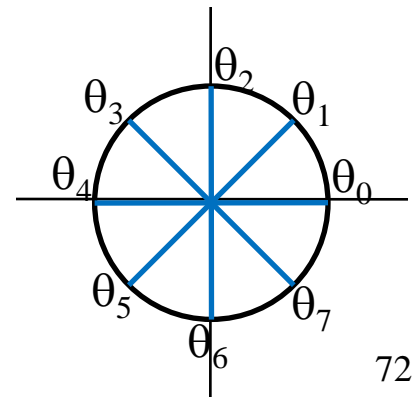


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# Moving Dot Products (Periodic Arrays)

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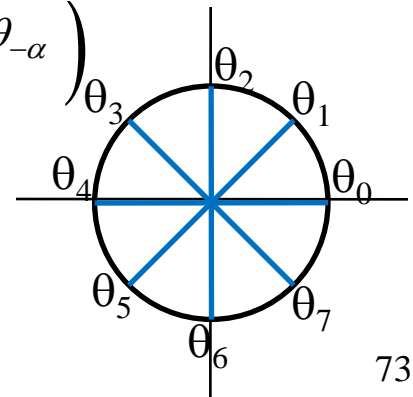
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So that:

$$\rho_\alpha(v_k[ ]) = (e^{ik\theta_0} \cdot e^{ik\theta_{-\alpha}}, \dots, e^{ik\theta_{n-1}} \cdot e^{ik\theta_{-\alpha}})$$





# Moving Dot Products (Periodic Arrays)

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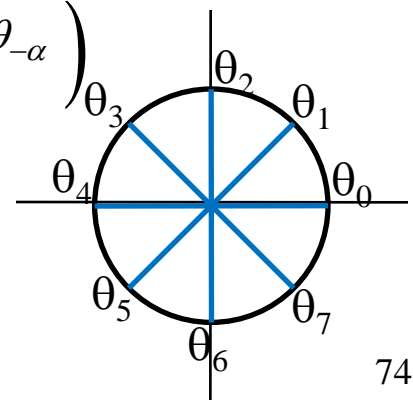
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So that:

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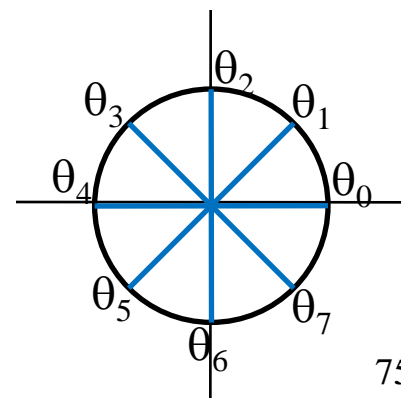


## Note 1

The periodic arrays:

$$v_k[ ] = (e^{ik\theta_0}, e^{ik\theta_1}, \dots, e^{ik\theta_{n-2}}, e^{ik\theta_{n-1}})$$

do not have unit norm!



# Moving Dot Products (Periodic Arrays)



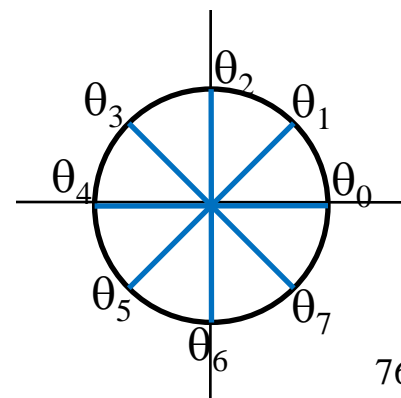
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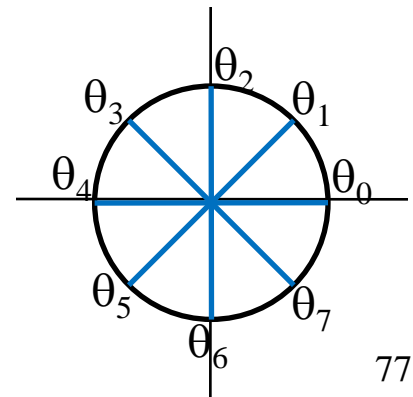
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# Moving Dot Products (Periodic Arrays)



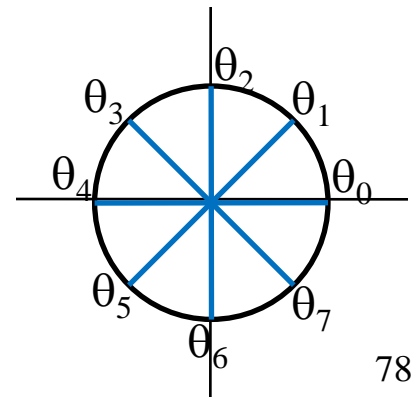
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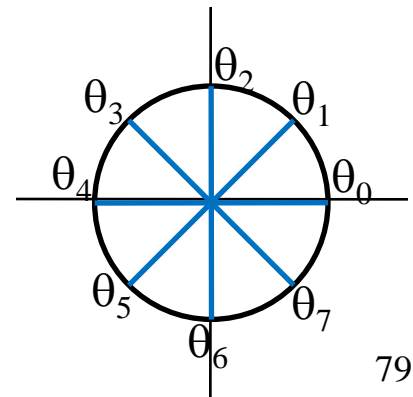
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# Moving Dot Products (Periodic Arrays)



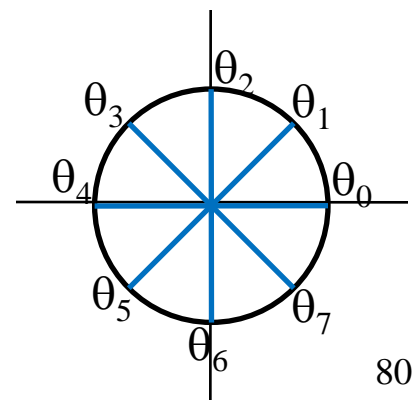
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# Moving Dot Products (Periodic Arrays)



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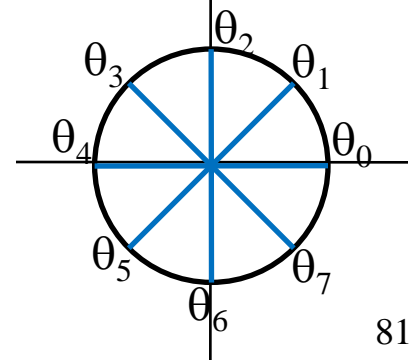
The periodic arrays:

$$v_k[ ] = (e^{ik\theta_0}, e^{ik\theta_1}, \dots, e^{ik\theta_{n-2}}, e^{ik\theta_{n-1}})$$

do not have unit norm!

We need to normalize these functions to make them unit-norm:

$$v_k = \sqrt{\frac{1}{n}} (e^{ik\theta_0}, e^{ik\theta_1}, \dots, e^{ik\theta_{n-2}}, e^{ik\theta_{n-1}})$$



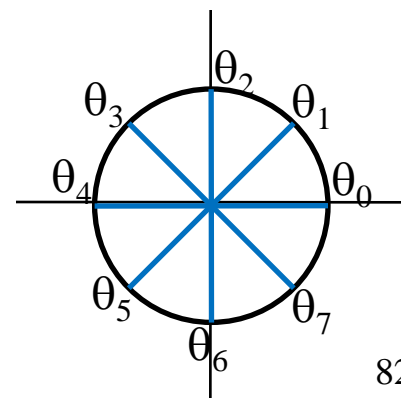
# Moving Dot Products (Periodic Arrays)



## Note 2

The arrays  $v_k[ ]$  and  $v_{k+n}[ ]$  are the same array:

$$\sqrt{n} \cdot v_{k+n}[ ] = \left( e^{i(k+n)\theta_0}, \dots, e^{i(k+n)\theta_{n-1}} \right)$$



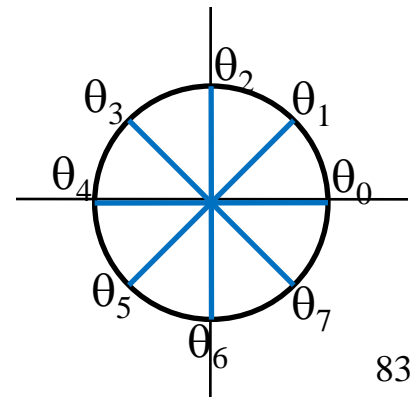
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# Moving Dot Products (Periodic Arrays)

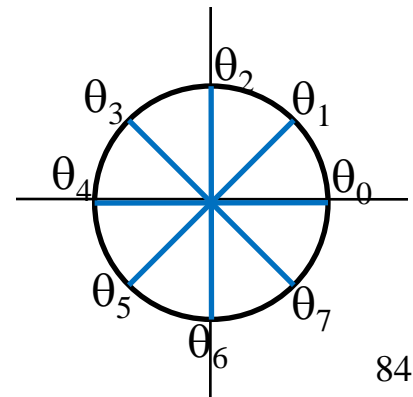
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But  $n\theta_j$  is just a multiple of  $2\pi$ :

$$n\theta_j = \frac{nj2\pi}{n} = j2\pi$$





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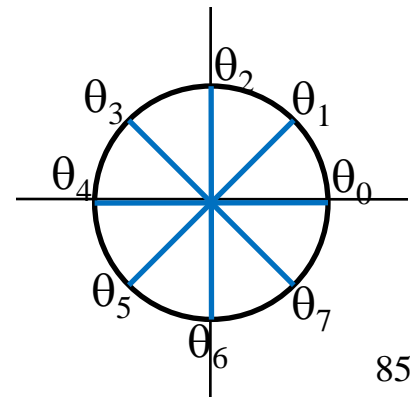
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$\updownarrow$

$$e^{in\theta_j} = 1$$





# Moving Dot Products (Periodic Arrays)

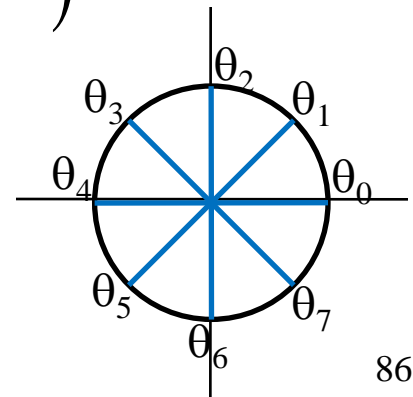
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But  $n\theta_j$  is just a multiple of  $2\pi$ , so

$$\begin{aligned}\sqrt{n} \cdot v_{k+n}[ ] &= \left( e^{ik\theta_0} \cdot e^{in\theta_0}, \dots, e^{ik\theta_{n-1}} \cdot e^{in\theta_{n-1}} \right) \\ &= \left( e^{ik\theta_0}, \dots, e^{ik\theta_{n-1}} \right) \\ &= \sqrt{n} \cdot v_k[ ]\end{aligned}$$

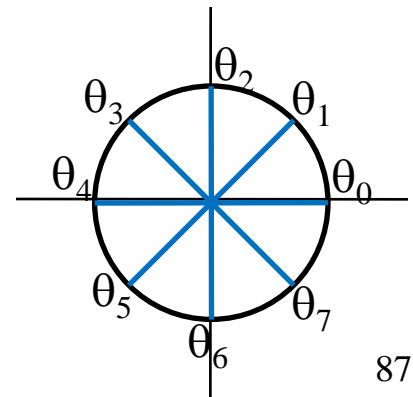


# Moving Dot Products (Periodic Arrays)



## Note 3

The arrays  $\{v_0[ ], \dots, v_{n-1}[ ]\}$  are linearly independent.



# Moving Dot Products (Periodic Arrays)

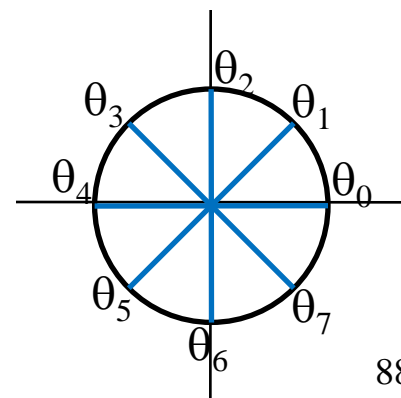


Thus, given two  $n$ -dimensional arrays,  $f[ ]$  and  $g[ ]$ , we can write:

$$f[ ] = \sum_{k=0}^{n-1} a_k v_k[ ] \quad \text{and} \quad g[ ] = \sum_{k=0}^{n-1} b_k v_k[ ]$$

to get:

$$(f[ ]^* g[ ])[\alpha] = \sum_{k=0}^{n-1} a_k \bar{b}_k \bar{\lambda}_k[\alpha]$$



# Moving Dot Products (Periodic Arrays)

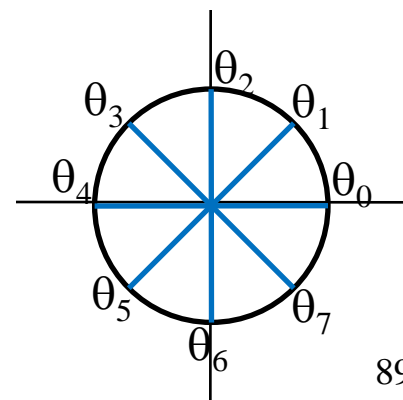


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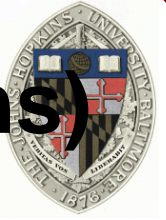


# Outline

## Review

### Moving Dot Products:

- One-Dimensional (Continuous)
- One-Dimensional (Discrete)
- **Higher-Dimensional**
- Computational Complexity



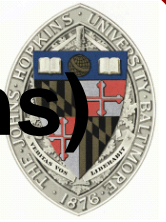
# Moving Dot Products (Higher Dimensions)

The same kind of method can be used for higher dimensions:

- Periodic functions in 2D

$$f_{lm}(\theta, \phi) = \sqrt{\frac{1}{(2\pi)^2}} e^{il\theta} \cdot e^{im\phi}$$

$$\bar{\lambda}_{lm}(\alpha, \beta) = \sqrt{(2\pi)^2} f_{lm}(\alpha, \beta)$$



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- Periodic functions in 3D

$$f_{lmn}(\theta, \phi, \varphi) = \sqrt{\frac{1}{(2\pi)^3}} e^{il\theta} \cdot e^{im\phi} \cdot e^{in\varphi}$$

$$\bar{\lambda}_{lmn}(\alpha, \beta, \gamma) = \sqrt{(2\pi)^3} f_{lmn}(\alpha, \beta, \gamma)$$



# Outline

## Review

## Moving Dot Products:

- One-Dimensional (Continuous)
- One-Dimensional (Discrete)
- Higher-Dimensional
- **Computational Complexity**



# Computational Complexity

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1. We need to express  $f[ ]$  and  $g[ ]$  in terms of the basis  $v_k[ ]$ :

$$f[ ] = \sum_{k=0}^{n-1} a_k v_k[ ] \quad \text{and} \quad g[ ] = \sum_{k=0}^{n-1} b_k v_k[ ]$$



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2. We need to multiply the coefficients:

$$(f[ ] * g[ ])[ ] = \sqrt{n} \sum_{k=0}^{n-1} a_k \bar{b}_k v_k[ ]$$



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# Computational Complexity

What do we need to do in order to compute the moving dot-product of two periodic,  $n$ -dimensional arrays  $f[ ]$  and  $g[ ]$ ?

The first and third steps are a change of bases.

This amounts to a matrix multiplication and can be as bad as quadratic in the dimension of the array.



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# Computational Complexity

What do we need to do in order to compute the moving dot-product of two periodic,  $n$ -dimensional arrays  $f[ ]$  and  $g[ ]$ ?

The Fast Fourier Transform (FFT) is an algorithm for expressing an array represented by samples at  $\{\theta_0, \dots, \theta_{n-1}\}$  as a linear sum of the  $v_k$ .

The Fast Inverse Fourier Transform (IFFT) is an algorithm for expressing an array represented as a linear sum of the  $v_k$  by samples at  $\{\theta_0, \dots, \theta_{n-1}\}$ .

Both take  $O(n \log n)$  time.



# Computational Complexity

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2. We need to multiply the coefficients:

$$(f[ ] * g[ ])[ ] = \sqrt{n} \sum_{k=0}^{n-1} a_k \bar{b}_k v_k[ ] \quad \boxed{O(n)}$$

3. We need to evaluate the moving dot-product at every index  $\alpha$ :

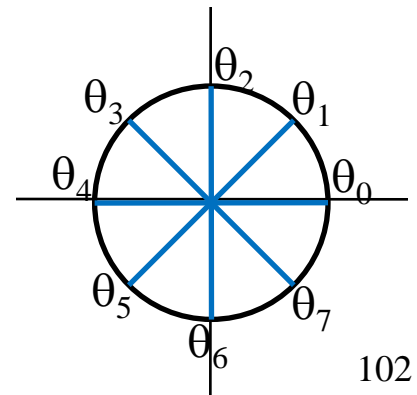
$$(f[ ] * g[ ])[\alpha] = \sqrt{n} \sum_{k=0}^{n-1} a_k \bar{b}_k v_k[\alpha] \quad \boxed{O(n \log n)}$$



# The Inverse Fourier Transform

The Fourier Transform is a change of basis transformation:

Evaluation Basis		Complex Exponential Basis
$(1, 0, \dots, 0, 0)$	$\xrightarrow{\text{Fourier Transform}}$	$(e^{i0\theta_0}, e^{i0\theta_1}, \dots, e^{i0\theta_{n-2}}, e^{i0\theta_{n-1}}) / \sqrt{n}$
$(0, 1, \dots, 0, 0)$		$(e^{i1\theta_0}, e^{i1\theta_1}, \dots, e^{i1\theta_{n-2}}, e^{i1\theta_{n-1}}) / \sqrt{n}$
$\vdots$		$\vdots$
$(0, 0, \dots, 1, 0)$		$(e^{i(n-2)\theta_0}, e^{i(n-2)\theta_1}, \dots, e^{i(n-2)\theta_{n-2}}, e^{i(n-2)\theta_{n-1}}) / \sqrt{n}$
$(0, 0, \dots, 0, 1)$		$(e^{i(n-1)\theta_0}, e^{i(n-1)\theta_1}, \dots, e^{i(n-1)\theta_{n-2}}, e^{i(n-1)\theta_{n-1}}) / \sqrt{n}$





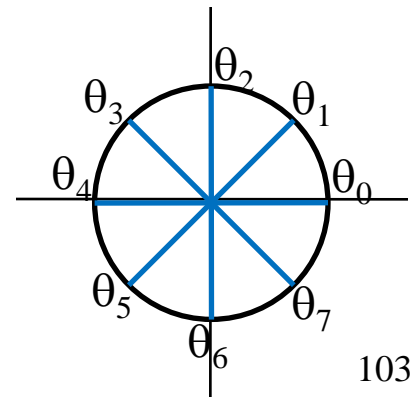
# The Inverse Fourier Transform

This can be represented by the matrix:

$$F = \sqrt{\frac{1}{n}} \begin{pmatrix} 1 & 1 & \dots & 1 & 1 \\ 1 & e^{i\theta} & \dots & e^{i(n-2)\theta} & e^{i(n-1)\theta} \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 1 & e^{i(n-2)\theta} & \dots & e^{i(n-2)(n-2)\theta} & e^{i(n-2)(n-1)\theta} \\ 1 & e^{i(n-1)\theta} & \dots & e^{i(n-2)(n-1)\theta} & e^{i(n-1)(n-1)\theta} \end{pmatrix}$$

where  $\theta$  is the angle:

$$\theta = \frac{2\pi}{n}$$

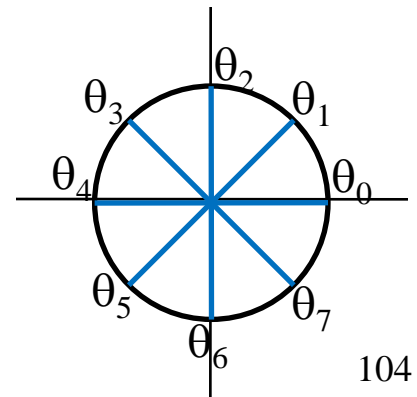




# The Inverse Fourier Transform

$$F = \sqrt{\frac{1}{n}} \begin{pmatrix} 1 & 1 & \dots & 1 & 1 \\ 1 & e^{i\theta} & \dots & e^{i(n-2)\theta} & e^{i(n-1)\theta} \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 1 & e^{i(n-2)\theta} & \dots & e^{i(n-2)(n-2)\theta} & e^{i(n-2)(n-1)\theta} \\ 1 & e^{i(n-1)\theta} & \dots & e^{i(n-2)(n-1)\theta} & e^{i(n-1)(n-1)\theta} \end{pmatrix}$$

Since both bases are orthogonal, the matrix is unitary, and the inverse Fourier transform is just the transpose conjugate of the forward Fourier transform.





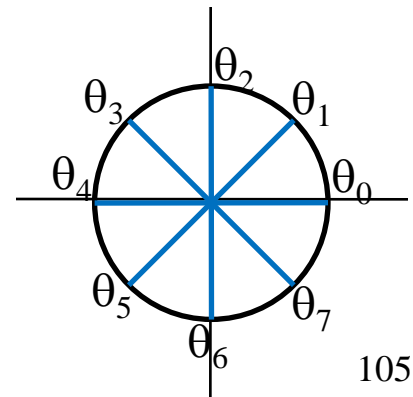
# The Inverse Fourier Transform

$$F = \sqrt{\frac{1}{n}} \begin{pmatrix} 1 & 1 & \dots & 1 & 1 \\ 1 & e^{i\theta} & \dots & e^{i(n-2)\theta} & e^{i(n-1)\theta} \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 1 & e^{i(n-2)\theta} & \dots & e^{i(n-2)(n-2)\theta} & e^{i(n-2)(n-1)\theta} \\ 1 & e^{i(n-1)\theta} & \dots & e^{i(n-2)(n-1)\theta} & e^{i(n-1)(n-1)\theta} \end{pmatrix}$$

In particular, given the Fourier coefficients:  
 $(a_0, \dots, a_{n-1})$

the inverse Fourier transform gives:

$$\overline{F}^t \begin{pmatrix} a_0 \\ \vdots \\ a_{n-1} \end{pmatrix}$$



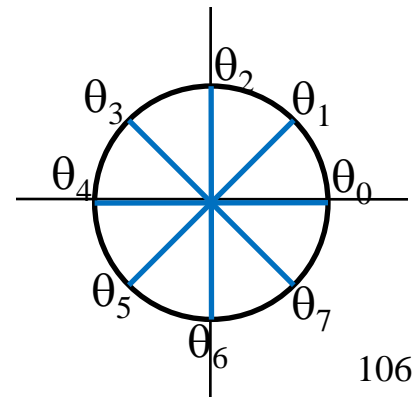


# The Inverse Fourier Transform

$$F = \sqrt{\frac{1}{n}} \begin{pmatrix} 1 & 1 & \dots & 1 & 1 \\ 1 & e^{i\theta} & \dots & e^{i(n-2)\theta} & e^{i(n-1)\theta} \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 1 & e^{i(n-2)\theta} & \dots & e^{i(n-2)(n-2)\theta} & e^{i(n-2)(n-1)\theta} \\ 1 & e^{i(n-1)\theta} & \dots & e^{i(n-2)(n-1)\theta} & e^{i(n-1)(n-1)\theta} \end{pmatrix}$$

Taking the double conjugate, we get:

$$\overline{\overline{F}}^t \begin{pmatrix} a_0 \\ \vdots \\ a_{n-1} \end{pmatrix} = \overline{\overline{F}}^t \begin{pmatrix} a_0 \\ \vdots \\ a_{n-1} \end{pmatrix}$$



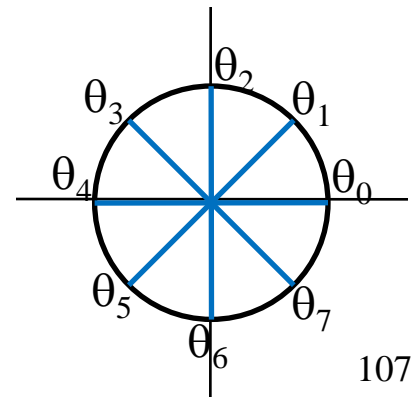


# The Inverse Fourier Transform

$$F = \sqrt{\frac{1}{n}} \begin{pmatrix} 1 & 1 & \dots & 1 & 1 \\ 1 & e^{i\theta} & \dots & e^{i(n-2)\theta} & e^{i(n-1)\theta} \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 1 & e^{i(n-2)\theta} & \dots & e^{i(n-2)(n-2)\theta} & e^{i(n-2)(n-1)\theta} \\ 1 & e^{i(n-1)\theta} & \dots & e^{i(n-2)(n-1)\theta} & e^{i(n-1)(n-1)\theta} \end{pmatrix}$$

Taking the double conjugate, we get:

$$\begin{aligned} \overline{\overline{F}}^t \begin{pmatrix} a_0 \\ \vdots \\ a_{n-1} \end{pmatrix} &= \overline{\overline{F}}^t \begin{pmatrix} a_0 \\ \vdots \\ a_{n-1} \end{pmatrix} \\ &= F^t \begin{pmatrix} \overline{a_0} \\ \vdots \\ \overline{a_{n-1}} \end{pmatrix} \end{aligned}$$



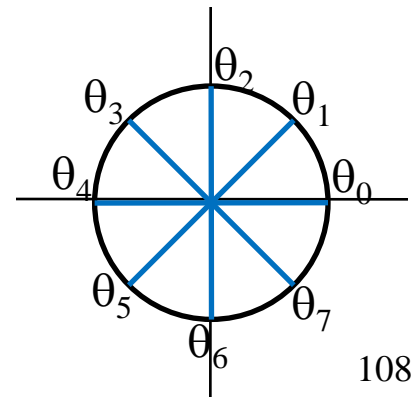


# The Inverse Fourier Transform

$$F = \sqrt{\frac{1}{n}} \begin{pmatrix} 1 & 1 & \dots & 1 & 1 \\ 1 & e^{i\theta} & \dots & e^{i(n-2)\theta} & e^{i(n-1)\theta} \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 1 & e^{i(n-2)\theta} & \dots & e^{i(n-2)(n-2)\theta} & e^{i(n-2)(n-1)\theta} \\ 1 & e^{i(n-1)\theta} & \dots & e^{i(n-2)(n-1)\theta} & e^{i(n-1)(n-1)\theta} \end{pmatrix}$$

Since  $F=F^t$ , this implies that:

$$\overline{F}^t \begin{pmatrix} a_0 \\ \vdots \\ a_{n-1} \end{pmatrix} = F \begin{pmatrix} \overline{a_0} \\ \vdots \\ \overline{a_{n-1}} \end{pmatrix}$$





# The Inverse Fourier Transform

$$\overline{F}^t \begin{pmatrix} a_0 \\ \vdots \\ a_{n-1} \end{pmatrix} = \overline{F \begin{pmatrix} \bar{a}_0 \\ \vdots \\ \bar{a}_{n-1} \end{pmatrix}}$$

Thus, we can compute the inverse Fourier transform by:

1. Taking the conjugate of the Fourier coefficients
2. Computing the forward Fourier transform
3. Taking the conjugate of the resultant coefficients.